# A Note About the Torsion of Null Curves in the 3-Dimensional Minkowski Spacetime and the Schwarzian Derivative 

Zbigniew Olszak ${ }^{\text {a }}$<br>Dedicated to the memory of Professor Włodzimierz Waliszewski

${ }^{a}$ Wroctazw University of Technology, Institute of Mathematics and Computer Science, 50-370 Wroctazw, Poland


#### Abstract

The main topic of this paper is to show that in the 3-dimensional Minkowski spacetime, the torsion of a null curve is equal to the Schwarzian derivative of a certain function appearing in a description of the curve. As applications, we obtain descriptions of the slant helices, and null curves for which the torsion is of the form $\tau=-2 \lambda s, s$ being the pseudo-arc parameter and $\lambda=$ const $\neq 0$.


## 1. Introduction

There are very many papers about geometric properties of null curves in the Minkowski spacetimes. We refer the monographs [4, 5], and the survey articles [3, 11, 12], etc.

On the other hand, there is the classical notion of the Schwarzian derivative in mathematical analysis. This notion has many important applications in mathematical analysis (real and complex) and differential geometry; see [6, 7, 13-15], etc. The author is specially inspired by the paper [7], where it is shown a strict relation between the Schwarzian derivative and the curvature of worldlines in 2-dimenional Lorentzian manifolds of constant curvature.

In the presented short paper, we will show that the torsion of a null curve in the 3-dimensional Minkowski spacetime $\mathbb{E}_{1}^{3}$ is equal to the Schwarzian derivative of a certain function appearing in a description of the curve. Descriptions of the slant helices are obtained, and null curves for which the torsion is given by $\tau=-2 \lambda s, s$ being the pseudo-arc parameter and $\lambda=$ const $\neq 0$.

## 2. Preliminaries

Let $\mathbb{E}_{1}^{3}$ be the 3-dimensional Minkowski spacetime, that is, the Cartesian $\mathbb{R}^{3}$ endowed with the standard Minkowski metric $g$ given with respect to the Cartesian coordinates $(x, y, z)$ by

$$
\begin{equation*}
g=d x \otimes d x+d y \otimes d y-d z \otimes d z \tag{1}
\end{equation*}
$$

or as the symmetric 2-form $g=d x^{2}+d y^{2}-d z^{2}$.

[^0]Let $\alpha: I \rightarrow \mathbb{E}_{1}^{3}$ be a null (ligth-like) curve in $\mathbb{E}_{1}^{3}, I$ being an open interval. Thus, $g\left(\alpha^{\prime}, \alpha^{\prime}\right)=0$, that is, $g\left(\alpha^{\prime}(t), \alpha^{\prime}(t)\right)=0$ for any $t \in I$. We also assume that the curve is non-degenerate, in the sense the three vector fields $\alpha^{\prime}, \alpha^{\prime \prime}, \alpha^{\prime \prime \prime}$ are linearly independent at every point of the curve.

Since $g\left(\alpha^{\prime}, \alpha^{\prime}\right)=0$ and $g\left(\alpha^{\prime}, \alpha^{\prime \prime}\right)=0$, it must be that $g\left(\alpha^{\prime \prime}, \alpha^{\prime \prime}\right)>0$. A parametrization of the null curve is said to be pseudo-arc (or distinguished) if $g\left(\alpha^{\prime \prime}, \alpha^{\prime \prime}\right)=1$. A null curve can always be parametrized by a pseudo-arc parameter. However, such a parameter is not uniquely defined. Precisely, for a null curve $\alpha$, if $s_{1}$ is a pseudo-arc parameter, then $s_{2}$ is a pseudo-arc parameter if and only if there exists a constant $c$ such that $s_{2}= \pm s_{1}+c$.

In the sequel, we assume that the parametrization of a null curve is pseudo-arc, and we denote such a parameter by $s$.

In the next section, we need the standard theorms concerning of null curves which can be formulated in the following manner (see e.g. [3, 5, 11, 12]):

Let $\alpha$ be a null curve in the 3-dimensional Minkowski spacetime $\mathbb{E}_{1}^{3}$. Then, there exists the only one Cartan moving frame $\left(\mathbf{L}=\alpha^{\prime}, \mathbf{N}, \mathbf{W}\right)$ and the function $\tau$ defined along the curve $\alpha$ and such that

$$
\begin{equation*}
g(\mathbf{L}, \mathbf{N})=g(\mathbf{W}, \mathbf{W})=1, \quad g(\mathbf{L}, \mathbf{L})=g(\mathbf{L}, \mathbf{W})=g(\mathbf{N}, \mathbf{N})=g(\mathbf{N}, \mathbf{W})=0 \tag{2}
\end{equation*}
$$

and the following system of differential equations

$$
\begin{equation*}
\mathbf{L}^{\prime}=\mathbf{W}, \quad \mathbf{N}^{\prime}=\tau \mathbf{W}, \quad \mathbf{W}^{\prime}=-\tau \mathbf{L}-\mathbf{N} \tag{3}
\end{equation*}
$$

is satisfied. These vector fileds are given by

$$
\begin{equation*}
\mathbf{L}=\alpha^{\prime}, \quad \mathbf{W}=\alpha^{\prime \prime}, \quad \mathbf{N}=-\alpha^{\prime \prime \prime}-\frac{1}{2} g\left(\alpha^{\prime \prime \prime}, \alpha^{\prime \prime \prime}\right) \alpha^{\prime} \tag{4}
\end{equation*}
$$

and the function $\tau$ by

$$
\begin{equation*}
\tau=\frac{1}{2} g\left(\alpha^{\prime \prime \prime}, \alpha^{\prime \prime \prime}\right) \tag{5}
\end{equation*}
$$

From these results it can be deduced that a given function $\tau$ on an open interval $I$, there exists the only one null curve $\alpha: I \rightarrow \mathbb{E}_{1}^{3}$ realizing (2) and (3) up to the orientation of this curve and up to the isometries of the Minkowski space $\mathbb{E}_{1}^{3}$.

The triple ( $\mathbf{L}, \mathbf{N}, \mathbf{W}$ ) defined in (4) is called the Frenet frame, the function $\tau$ defined in (5) is called the torsion, and the equations (3) are called the Frenet equations of the null curve $\alpha$. Since

$$
\begin{equation*}
\operatorname{det}[\mathbf{L}, \mathbf{N}, \mathbf{W}]=\operatorname{det}\left[\alpha^{\prime}, \alpha^{\prime \prime}, \alpha^{\prime \prime \prime}\right] \tag{6}
\end{equation*}
$$

the frames $(\mathbf{L}, \mathbf{N}, \mathbf{W})$ and $\left(\alpha^{\prime}, \alpha^{\prime \prime}, \alpha^{\prime \prime \prime}\right)$ have the same orientations.
In the following section, we are going to expresse the torsion $\tau$ and the frame ( $\mathbf{L}, \mathbf{N}, \mathbf{W}$ ) with the help of a special function related to a pseudo-arc parametrization of a null curve in $\mathbf{E}_{1}^{3}$.

## 3. A description of the torsion

Let $\alpha: I \rightarrow \mathbb{E}_{1}^{3}$ be a null curve. Simplifying denotations, we write $\alpha(s)=(x(s), y(s), z(s)), s \in I$, where $s$ is a pseudo-arc parameter, and $x(s), y(s), z(s)$ are certain functions of $s$. Then, we have

$$
\alpha^{\prime}=\left.x^{\prime} \frac{\partial}{\partial x}\right|_{\alpha}+\left.y^{\prime} \frac{\partial}{\partial y}\right|_{\alpha}+\left.z^{\prime} \frac{\partial}{\partial z}\right|_{\alpha} .
$$

For simplicity, instead of that, we will write $\alpha^{\prime}=\left(x^{\prime}, y^{\prime}, z^{\prime}\right)$. And in the similar manner, the next derivatives of $\alpha$ will be written, e.g., $\alpha^{\prime \prime}=\left(x^{\prime \prime}, y^{\prime \prime}, z^{\prime \prime}\right)$.

Using (1), our two assumptions: $g\left(\alpha^{\prime}, \alpha^{\prime}\right)=0$ (the nullity condition), and $g\left(\alpha^{\prime \prime}, \alpha^{\prime \prime}\right)=1$ (the pseudo-arc parametrization) give the following two equalities

$$
\begin{array}{r}
x^{\prime 2}+y^{\prime 2}-z^{\prime 2}=0 \\
x^{\prime \prime 2}+y^{\prime \prime 2}-z^{\prime \prime 2}=1 \tag{8}
\end{array}
$$

One notes that the shapes of the equalities (7) and (8) exclude the situation when at least one of the functions $x^{\prime}, y^{\prime}, z^{\prime}$ vanishes on an open subinterval of $I$. In the sequel, restricting slightly the assumptions, we will consider only the case when $x^{\prime} \neq 0, y^{\prime} \neq 0$ and $z^{\prime} \neq 0$ on $I$.

It is a standard and elementary idea that from (7), it follows that

$$
\begin{equation*}
x^{\prime}=h, \quad y^{\prime}=\frac{h}{2}\left(f-\frac{1}{f}\right), \quad z^{\prime}=\frac{h}{2}\left(f+\frac{1}{f}\right) \tag{9}
\end{equation*}
$$

$f$ and $h$ being certain non-zero functions on $I$. Hence,

$$
\begin{aligned}
& x^{\prime \prime}=h^{\prime} \\
& y^{\prime \prime}=\frac{f h^{\prime}\left(f^{2}-1\right)+h f^{\prime}\left(f^{2}+1\right)}{2 f^{2}} \\
& z^{\prime \prime}=\frac{f h^{\prime}\left(f^{2}+1\right)+h f^{\prime}\left(f^{2}-1\right)}{2 f^{2}} .
\end{aligned}
$$

In view of the above relations, the equality (8) turns into $h^{2} f^{\prime 2}=f^{2}$. Hence, $f^{\prime}$ is non-zero (and has constant sign) on I. Consequently,

$$
h=\varepsilon \frac{f}{f^{\prime}}, \varepsilon= \pm 1
$$

Thus, for the vector field $\mathbf{L}$ (cf. (4)), we have

$$
\begin{equation*}
\mathbf{L}=\alpha^{\prime}=\frac{\varepsilon}{2 f^{\prime}}\left(2 f, f^{2}-1, f^{2}+1\right) \tag{10}
\end{equation*}
$$

Consequently, we get the following description of the curve $\alpha$

$$
\alpha(s)=\alpha\left(s_{0}\right)+\frac{\varepsilon}{2} \int_{s_{0}}^{s} \frac{1}{f^{\prime}(t)}\left(2 f(t), f^{2}(t)-1, f^{2}(t)+1\right) d t, s, s_{0} \in I
$$

Conversely, if a curve $\alpha$ is given by the last formula, then (7) and (8) are fulfilled so that the curve is null and not geodesic, and the parameter $s$ is distinguish.

From (10), we obtain for the vector field $\mathbf{W}$ (cf. (4)),

$$
\begin{equation*}
\mathbf{W}=\alpha^{\prime \prime}=-\frac{\varepsilon f^{\prime \prime}}{2 f^{\prime 2}}\left(2 f, f^{2}-1, f^{2}+1\right)+\varepsilon(1, f, f) \tag{11}
\end{equation*}
$$

From (11), we find

$$
\begin{equation*}
\alpha^{\prime \prime \prime}=\varepsilon \frac{2 f^{\prime \prime 2}-f^{\prime} f^{\prime \prime \prime}}{2 f^{\prime 3}}\left(2 f, f^{2}-1, f^{2}+1\right)-\frac{\varepsilon f^{\prime \prime}}{f^{\prime}}(1, f, f)+\varepsilon f^{\prime}(0,1,1) \tag{12}
\end{equation*}
$$

To compute $g\left(\alpha^{\prime \prime \prime}, \alpha^{\prime \prime \prime}\right)$, using (1), we find at first the following

$$
\begin{aligned}
& g\left(\left(2 f, f^{2}-1, f^{2}+1\right),\left(2 f, f^{2}-1, f^{2}+1\right)\right)=0 \\
& g\left(\left(2 f, f^{2}-1, f^{2}+1\right),(1, f, f)\right)=0, \quad g\left(\left(2 f, f^{2}-1, f^{2}+1\right),(0,1,1)\right)=-2 \\
& g((1, f, f),(1, f, f))=1, \quad g((1, f, f),(0,1,1))=0, \quad g((0,1,1),(0,1,1))=0 .
\end{aligned}
$$

Then, having (12) and applying the above formulas, we get

$$
\begin{equation*}
g\left(\alpha^{\prime \prime \prime}, \alpha^{\prime \prime \prime}\right)=\frac{2 f^{\prime} f^{\prime \prime \prime}-3 f^{\prime \prime 2}}{f^{\prime 2}} \tag{13}
\end{equation*}
$$

In view of (13) and (5), the torsion must be of the form

$$
\begin{equation*}
\tau=\frac{2 f^{\prime} f^{\prime \prime \prime}-3 f^{\prime \prime \prime}}{2 f^{\prime 2}}=\left(\frac{f^{\prime \prime}}{f^{\prime}}\right)^{\prime}-\frac{1}{2}\left(\frac{f^{\prime \prime}}{f^{\prime}}\right)^{2} \tag{14}
\end{equation*}
$$

Now, it is important to note that the right hand side of the formula (14) is just the Schwarzian derivative of the function $f$, which is usually denoted by $S(f)$. Thus, $\tau=S(f)$.

Finally, applying (10), (12) and (13) into (4), we find the vector field

$$
\begin{equation*}
\mathbf{N}=-\frac{\varepsilon f^{\prime \prime 2}}{4 f^{\prime 3}}\left(2 f, f^{2}-1, f^{2}+1\right)+\frac{\varepsilon f^{\prime \prime}}{f^{\prime}}(1, f, f)-\varepsilon f^{\prime}(0,1,1) \tag{15}
\end{equation*}
$$

Summarizing the above considerations, we can formulate the following theorem.
Theorem 1. Let $\mathbb{E}_{1}^{3}$ be the 3-dimensional Minkowski spacetime. Any (non-degenerate) null curve $\alpha$ in $\mathbb{E}_{1}^{3}$ can be parametrized in the following way

$$
\begin{equation*}
\alpha(s)=\alpha\left(s_{0}\right)+\frac{\varepsilon}{2} \int_{s_{0}}^{s} \frac{1}{f^{\prime}(t)}\left(2 f(t), f^{2}(t)-1, f^{2}(t)+1\right) d t, s, s_{0} \in I \tag{16}
\end{equation*}
$$

where s is a pseudo-arc parameter, I is a certain open interval, $f$ is a non-zero function with non-zero derivative $f^{\prime}$ on I. The torsion $\tau$ of such a curve is equal to the Schwarzian derivative of the function $f$, that is,

$$
\begin{equation*}
\tau=S(f)=\left(\frac{f^{\prime \prime}}{f^{\prime}}\right)^{\prime}-\frac{1}{2}\left(\frac{f^{\prime \prime}}{f^{\prime}}\right)^{2} \tag{17}
\end{equation*}
$$

The vector fields forming the Frenet frame of the curve $\alpha$ are given by the formulas (10), (11) and (15).
Remark 1. Applying formulas (10), (11), (15), it can be verified that

$$
\operatorname{det}[\mathbf{L}, \mathbf{N}, \mathbf{W}]=\varepsilon
$$

This together with (6) implies that the constant $\varepsilon$ appering in (16) corresponds to the orientation of the curve $\alpha$. Note that the torsion does not depend on the orientation of the curve. Moreover, the torsion and the orientation does not depend on the sign of the function $f$.

Remark 2. The Schwarzian derivative $S$ is an invariant of a fractional-linear transformation $T$ of the 1-dimensional real projective space $\mathbb{R} P^{1}=\mathbb{R} \cup \infty$ (cf.e.g. [13]). That is, $S(T \circ f)=S(f)$ if $f$ is a function on $\mathbb{R} P^{1}$ and $T: \mathbb{R} P^{1} \rightarrow \mathbb{R} P^{1}$ is given by

$$
\begin{equation*}
T(r)=\frac{a r+b}{c r+d}, r \in \mathbb{R} P^{1}, a, b, c, d \in \mathbb{R}, a d-b c \neq 0 \tag{18}
\end{equation*}
$$

We can apply the above fact seeking for null curves with given torsion $\tau$. However, we should be careful since the domains of our functions $f$ and $T \circ f$ may be defined only on some open subintervals lying on the real line $\mathbb{R}$.

## 4. Null Cartan helices

It is well-known that there are exactly three types of null curves with constant torsion in the Minkowski spacetime $\mathbb{E}_{1}^{3}$ (cf e.g., [9]) up to the orientation of the curve and up to the isometries of the space. They are often called the null Cartan helices.

As a first application of the results from the previous section, we demonstrate how these classes of curves can be recovered from their torsions.
(a) For $f(s)=s$, it holds $S(f)=0$. In (16), we put $f(s)=s, s_{0}=0, \alpha\left(s_{0}\right)=(0,0,0), \varepsilon=1$. Then, we obtain the curve

$$
\alpha(s)=\frac{1}{6}\left(3 s^{2}, s^{3}-3 s, s^{3}+3 s\right)
$$

for which by (17) we have $\tau=0$. Thus, $\alpha$ is a positively oriented null Cartan helix of zero torsion.
(b) For $f(s)=-\cot (c s / 2)$, it holds $S(f)=c^{2} / 2$. In (16), we put

$$
f(s)=-\cot \frac{c s}{2}, \alpha(0)=\left(\frac{1}{c^{2}}, 0,0\right), \varepsilon=1, c=\text { const. }>0
$$

Then, we obtain the curve

$$
\alpha(s)=\frac{1}{c^{2}}(\cos (c s), \sin (c s), c s),
$$

for which by (17) it holds $\tau=c^{2} / 2$. Thus, $\alpha$ is a positively oriented null Cartan helix of constant positive torsion.
(c) For $f(s)=e^{c s}$, it holds $S(f)=-c^{2} / 2$. In (16), we put

$$
f(s)=e^{c s}, \alpha(0)=\left(0, \frac{1}{c^{2}}, 0\right), \varepsilon=1, c=\text { const } .>0
$$

Then, we obtain the curve

$$
\alpha(s)=\frac{1}{c^{2}}(c s, \cosh (c s), \sinh (c s)),
$$

for which by (17) we have $\tau=-c^{2} / 2$. Thus, $\alpha$ is a positively oriented null Cartan helix of constant negative torsion.

Thus, we have seen the following:
Corollary 1. Null helices in $\mathbb{E}_{1}^{3}$ form the three classes described in (a) - (c) in the above. The description is valid up to the pseudo-arc parameter changies, up to the orientation of the curve, and up to the isometries of the space.

A curve $\alpha: I \rightarrow \mathbb{E}_{1}^{3}$ is called a general (or generalized) helix if there exists a non-zero vector $V$ in $\mathbb{E}_{1}^{3}$ such that $g\left(\alpha^{\prime}, V\right)=$ const.; cf. [8, 9, 17], etc. This means that tangent indicatrix is laid in a plane or, equivalently, there exists a non-zero constant vector $V$ in $\mathbb{E}_{1}^{3}$ for which $g\left(\alpha^{\prime \prime}, V\right)=0$, that is, $V$ is orthogonal to the acceleration vector field $\alpha^{\prime \prime}$.

For null curves, it is already proved that null general helices in $\mathbb{E}_{1}^{3}$ are precisely the null Cartan helices; cf. ibidem.

## 5. Null slant helices

Following the ideas of $[1,2,10]$, a slant helix is defined to be the curve (null as well as non-null) in $\mathbb{E}_{1}^{3}$ which satisfies the condition

$$
\begin{equation*}
g\left(\alpha^{\prime \prime}, V\right)=c=\text { const } \tag{19}
\end{equation*}
$$

along the curve $\alpha$, where $V$ is a constant vector. Thus, a general helix is a slant helix with $c=0$. Conversely, a slant helix with $c=0$ becomes a general helix. In [1, Theorem 1.4], it is proved that a null curve in $\mathbb{E}_{1}^{3}$ is a slant helix if and only if its torsion is given by

$$
\begin{equation*}
\tau=\frac{a}{(c s+b)^{2}}, a, b, c=\text { const. } \tag{20}
\end{equation*}
$$

where $c$ is just the constant realizing (19).
As the second applications of the results from Section 3, we will describe the null slant helices in $\mathbb{E}_{1}^{3}$ which are different from the usual helices ( $a \neq 0$ and $c \neq 0$ in (20)).

Note that moving the pseudo-arc parameter $s$ into $s-b / c$ and next modifying slightly the constant $a$, we can write the condition (20) as

$$
\begin{equation*}
\tau=\frac{a}{2 s^{2}}, a=\mathrm{const} \neq 0 \tag{21}
\end{equation*}
$$

We can also assume that $s>0$. Using (2) and (3), it can be checked that when the relation (21) is fulfilled, then for the vector

$$
V=-\frac{a}{2 s} \mathbf{L}+s \mathbf{N}+\mathbf{W}
$$

it holds $V^{\prime}=0$ and $g\left(\alpha^{\prime \prime}, V\right)=g(\mathbf{W}, V)=1$ (cf. ibidem).
(a) In (16), we put

$$
f(s)=\ln s, s_{0}=1, \alpha\left(s_{0}\right)=\frac{1}{8}(-2,-1,3), \varepsilon=1 .
$$

Then, we obtain the curve

$$
\alpha(s)=\frac{s^{2}}{8}\left(2(2 \ln s-1), 2 \ln ^{2} s-2 \ln s-1,2 \ln ^{2} s-2 \ln s+3\right)
$$

for which by (17) it holds

$$
\tau=S(f)=\frac{1}{2 s^{2}}
$$

Thus, $\alpha$ is a slant helix realizing (21) with $a=1$.
(b) Let $a>1$ and $b=\sqrt{a-1}>0$. In (16), we put

$$
f(s)=\tan \left(\frac{1}{2} \ln s^{b}\right), s_{0}=1, \alpha\left(s_{0}\right)=\frac{1}{b}\left(-\frac{b}{b^{2}+4},-\frac{2}{b^{2}+4}, \frac{1}{2}\right), \varepsilon=1
$$

Then, we obtain the curve

$$
\alpha(s)=\frac{s^{2}}{b}\left(\frac{2 \sin \left(\ln s^{b}\right)-b \cos \left(\ln s^{b}\right)}{b^{2}+4},-\frac{2 \cos \left(\ln s^{b}\right)+b \sin \left(\ln s^{b}\right)}{b^{2}+4}, \frac{1}{2}\right)
$$

for which by (17) it holds

$$
\tau=S(f)=\frac{1+b^{2}}{2 s^{2}}=\frac{a}{2 s^{2}}
$$

Thus, $\alpha$ is a slant helix realizing (21) with $a>1$.
(c) Let $0 \neq a<1$. Then for $b=\sqrt{1-a}$, we have $b>0$ and $b \neq 1$. Consider the case $a \neq-3$, that is, $b \neq 2$. In (16), we put

$$
f(s)=s^{-b}, s_{0}=1, \alpha\left(s_{0}\right)=\frac{1}{2 b}\left(-1, \frac{2 b}{b^{2}-4}, \frac{4}{b^{2}-4}\right), \varepsilon=1 .
$$

Then, we obtain the curve

$$
\alpha(s)=\frac{s^{2}}{2 b}\left(-1, \frac{s^{-b}}{b-2}+\frac{s^{b}}{b+2}, \frac{s^{-b}}{b-2}-\frac{s^{b}}{b+2}\right)
$$

for which by (17) it holds

$$
\tau=S(f)=\frac{1-b^{2}}{2 s^{2}}=\frac{a}{2 s^{2}}
$$

Thus, $\alpha$ is a slant helix realizing (21) with $-3 \neq a<1$.
(d) In (16), we put

$$
f(s)=\frac{1}{s^{2}}, s_{0}=1, \alpha\left(s_{0}\right)=\frac{1}{16}(-4,1,-1), \varepsilon=1
$$

Then, we obtain the curve

$$
\alpha(s)=\frac{1}{16}\left(-4 s^{2}, s^{4}-4 \ln s,-s^{4}-4 \ln s\right)
$$

for which by (17) it holds

$$
\tau=S(f)=-\frac{3}{2 s^{2}}
$$

Thus, $\alpha$ is a slant helix realizing (21) with $a=-3$.
Thus, we have shown the following:
Corollary 2. Null slant helices in $\mathbb{E}_{1}^{3}$ form the four classes described in $(a)-(d)$ in the above. The description is valid up to the pseudo-arc parameter changies, up to the orientation of the curve, and up to the isometries of the space.

## 6. Null curves with the torsion proportional to the pseudo-arc parameter

In this section, we determine the null curves in $\mathbb{E}_{1}^{3}$ for which $\tau=-2 \lambda s, \lambda=$ const. $\neq 0$. We will use the formula (17).

According to our Theorem, we need at first to find a solution of the differential equation

$$
\begin{equation*}
\left(\frac{f^{\prime \prime}}{f^{\prime}}\right)^{\prime}-\frac{1}{2}\left(\frac{f^{\prime \prime}}{f^{\prime}}\right)^{2}=-2 \lambda s \tag{22}
\end{equation*}
$$

We seek for solutions of this equation in the form

$$
\begin{equation*}
f(s)=\int \frac{d s}{\phi^{2}(s)} \tag{23}
\end{equation*}
$$

$\phi$ being an unknown fucntion. Then the equation (22) becomes the following differential equation

$$
\begin{equation*}
\phi^{\prime \prime}-\lambda s \phi=0 \tag{24}
\end{equation*}
$$

The general solution of the above equation is

$$
\phi(s)=c_{1} \operatorname{Ai}(\mu s)+c_{2} \operatorname{Bi}(\mu s), \mu=\sqrt[3]{\lambda}, c_{1}, c_{2}=\text { const }
$$

where Ai and Bi are the Airy functions of the first and second kind, respectively. For the solutions of (24) and for the special Airy functions, we refer [16], [18], etc. In the below calculations, we use the basic properties of these functions.

For our purpose, we take the only one solution of (24), say $\phi(s)=\operatorname{Ai}(\mu s)$. Then, from (23) we get

$$
\begin{equation*}
f(s)=\frac{\pi}{\mu} \cdot \frac{\operatorname{Bi}(\mu s)}{\operatorname{Ai}(\mu s)} . \tag{25}
\end{equation*}
$$

Next,

$$
\begin{equation*}
f^{\prime}(s)=\frac{1}{\operatorname{Ai}^{2}(\mu s)} \tag{26}
\end{equation*}
$$

Having (10) with $\varepsilon=1$, and using (25) and (26), we can write $\alpha^{\prime}$ as

$$
\alpha^{\prime}(s)=\left(\frac{\pi}{\mu} \operatorname{Ai}(\mu s) \operatorname{Bi}(\mu s), \frac{1}{2 \mu^{2}}\left(\pi^{2} \operatorname{Bi}^{2}(\mu s)-\mu^{2} \operatorname{Ai}^{2}(\mu s)\right), \frac{1}{2 \mu^{2}}\left(\pi^{2} \operatorname{Bi}^{2}(\mu s)+\mu^{2} \operatorname{Ai}^{2}(\mu s)\right)\right)
$$

The integration of the last equality gives the following curve

$$
\begin{align*}
& \alpha(s)=\left(\frac{\pi}{\mu^{2}}\left(\mu s \operatorname{Ai}(\mu s) \operatorname{Bi}(\mu s)-\operatorname{Ai}^{\prime}(\mu s) \operatorname{Bi}^{\prime}(\mu s)\right),\right. \\
& \frac{1}{2 \mu^{3}}\left(\pi^{2}\left(\mu s \operatorname{Bi}^{2}(\mu s)-\operatorname{Bi}^{\prime 2}(\mu s)\right)-\mu^{3} s \operatorname{Ai}^{2}(\mu s)+\mu^{2} \operatorname{Ai}^{\prime 2}(\mu s)\right), \\
&\left.\frac{1}{2 \mu^{3}}\left(\pi^{2}\left(\mu s \operatorname{Bi}^{2}(\mu s)-\operatorname{Bi}^{\prime 2}(\mu s)\right)+\mu^{3} s \operatorname{Ai}^{2}(\mu s)-\mu^{2} \operatorname{Ai}^{\prime 2}(\mu s)\right)\right), \tag{27}
\end{align*}
$$

if the the initial condition at $s_{0}=0$ is

$$
\alpha(0)=\frac{1}{2 \sqrt[3]{9} \mu^{3} \Gamma^{2}\left(\frac{1}{3}\right)}\left(2 \sqrt{3} \mu \pi, \mu^{2}-3 \pi^{2},-\mu^{2}-3 \pi^{2}\right)
$$

Thus, we can formulate the following:
Corollary 3. Null curves in $\mathbb{E}_{1}^{3}$ for which $\tau=-2 \lambda s, \lambda=$ const. $\neq 0$, are given by the formula (27) with $\mu=\sqrt[3]{\lambda}$ up to the pseudo-arc parameter changies, up to the orientation of the curve, and up to the isometries of the space.

## References

[1] A. T. Ali and R. López, Slant helices in Minkowski space $\mathbb{E}_{1}^{3}$, Journal of the Korean Mathematical Society 48 (2011) $159-167$.
[2] J. H. Choi and Y. H. Kim, Note on null helices in $\mathbb{E}_{1}^{3}$, Bulletin of the Korean Mathematical Society 50 (2013) 885-899.
[3] K. L. Duggal, A report on canonical null curves and screen distributions for lightlike geometry, Acta Applicandae Mathematicae 95 (2007) 135-149.
[4] K. L. Duggal and A. Bejancu, Lightlike submanifolds of semi-Riemannian manifolds and applications, Kluwer Academic Publishers, Dordrecht, 1996.
[5] K. L. Duggal and D. H. Jin, Null curves and hypersurfaces of semi-Riemannian manifolds, World Scientific Publishing, Hackensack, 2007.
[6] C. Duval and L. Guieu, The Virasoro group and Lorentzian surfaces: the hyperboloid of one sheet, Journal of Geometry and Physics 33 (2000) 103-127.
[7] C. Duval and V. Ovsienko, Lorentzian worldlines and the Schwarzian derivative, Functional Analysis and its Applications 34 (2000), No. 2, 135-137; translation from Funktsionalnyj Analiz i Ego Prilozheniya 34 (2000), No. 2, 69-72.
[8] A. Ferrández, A. Giménez and P. Lucas, Null helices in Lorentzian space forms, International Journal of Modern Physics A 16 (2001) 4845-4863.
[9] A. Ferrández, A. Giménez and P. Lucas, Null generalized helices in Lorentz-Minkowski spaces, Journal of Physics A: Mathematical and General 35 (2002) 8243-8251.
[10] F. Gökçelik and İ. Gök, Null W-slant helices in $\mathbf{E}_{1}^{3}$, Journal of Mathematical Analysis and Applications 420 (2014) 222-241.
[11] J.-I. Inoguchi and S. Lee, Null curves in Minkowski 3-space, International Electronic Journal of Geometry 1 (2008), No. 2, 40-83.
[12] R. Lopez, Differential geometry of curves and surfaces in Lorentz-Minkowski space, International Electronic Journal of Geometry 7 (2014), No. 1, 44-107.
[13] B. Osgood, Old and new on the Schwarzian derivative, In: Duren, Peter (ed.) et al., Quasiconformal mappings and analysis. Proceedings of the international symposium, Ann Arbor, MI, USA, August 1995, pp. 275-308, New York, NY: Springer, 1998.
[14] B. Osgood and D. Stowe, The Schwarzian derivative and conformal mapping of Riemannian manifolds, Duke Mathematical Journal 67 (1992) 57-99.
[15] V. Ovsienko and S. Tabachnikov, What is the Schwarzian derivative?, Notices of the American Mathematical Society 56 (2009) 34-36.
[16] A. D. Polyanin and V. F. Zaitsev, Handbook of exact solutions for ordinary differential equations, 2nd Edition, Chapman \& Hall/CRC, Boca Raton, 2003.
[17] B. Şahin, E. Kiliç and R. Güneş, Null helices in $\mathbb{R}_{1}^{3}$, Differential Geometry - Dynamical Systems 3 (2001), No. 2, 31-36.
[18] O. Vallée and M. Soares, Airy functions and applications to physics, Imperial College Press, London, 2004.


[^0]:    2010 Mathematics Subject Classification. Primary 53A35; Secondary 53B30, 53B50, 53C50.
    Keywords. Minkowski spacetime of dimension 3; null curve; light-like curve; torsion of a curve; null helix; slant helix; Airy function; Schwarzian derivative.

    Received: 2 September 2014; Accepted: 21 November 2014
    Communicated by Ljubica Velimirović and Mića Stanković
    Email address: zbigniew.olszak@pwr.edu.pl (Zbigniew Olszak)

