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# Non-fundamental trunc-simplex tilings and their optimal hyperball packings and coverings in hyperbolic space I. For families F1-F4

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### To Honour of János Bolyai on the 220th Anniversary of His Birth

**Abstract.** Supergroups of some hyperbolic space groups are classified as a continuation of our former works. Fundamental domains will be integer parts of truncated tetrahedra belonging to families F1 - F4, for a while, by the notation of E. Molnár et al. in 2006. Our paper relies basically on this work. As an application, optimal congruent hyperball packings and coverings to the truncation base planes with their very good densities are computed. This covering density is better than the conjecture of L. Fejes Tóth for balls and horoballs in 1964.

# 1. Introduction

### 1.1. Short history

Hyperbolic space groups are isometry groups, acting discontinuously on hyperbolic 3-space  $\mathbb{H}^3$  with compact fundamental domains. It will be investigated some series of such groups by looking for their fundamental domains. Face pairing identifications on a given polyhedron give us generators and relations for a space group by the algorithmically generalized Poincaré Theorem [4], [6], [11] (as in subsection 1.2).

The simplest fundamental domains are 3-simplices (tetrahedra) and their integer parts by inner symmetries. In the process of classifying the fundamental simplices, 64 combinatorially different face pairings of fundamental simplices were determined [28], [8], and also 35 solid transitive non-fundamental simplex identifications [10]. I. K. Zhuk [28] classified Euclidean and hyperbolic fundamental simplices of finite volume. An algorithmic procedure was given by E. Molnár and I. Prok [7]. In [8–10] the authors summarized all those results, arranging identified simplices into 32 families. Each of them is characterized by the so-called maximal series of simplex tilings, i.e. maximal symmetry groups with smallest fundamental

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domains. The other members are derived by symmetry breaking of this maximal one. These principles are also discussed by *D*-symbols in [10], see also 1.2.

Some complete cases of supergroups with fundamental truncated simplices (shortly trunc-simplices) are discussed in [5, 14–19]. As a classification, supergroups of the groups with fundamental trunc-simplices are given in [20]. Especially, supergroups of the groups with fundamental simplices belonging to families F1-F4, by the notations given in [10], are investigated in [5, 15, 18, 19]. For that purpose, the stabilizer group of the corresponding vertex figure is analyzed. The new method, for such analysis, established e.g. in the summary works [19, 20], is also confirmed by the results given in [3].

We think that our systematic investigations on hyperbolic space groups have newer applications in todays material sciences (e.g. fullerenes, nanomaterials). Extremal problems related to them (e.g. optimal ball packing and covering, see [13] select these infinite series to choose models for optimal real materials.

Thus, the discovery of János Bolyai and Nikolai Ivanovich Lobachevsky gets new and new importances and perspectives.

### 1.2. Preliminaries

Generators and relations for a space group *G* with a given polyhedron *P* as fundamental domain can be obtained by the Poincarè theorem, applying Poincarè algorithm. In this paper such polyhedra will be simplices and trunc-simplices, for overview see e.g. [19, 20], and here mainly their integer parts.

It is necessary to consider all *face pairing identifications* of such domains. These are isometries which generate a group *G* and induce subdivision of faces and directed edge segments of *P* into equivalence classes, such that an edge segment does not contain two *G*-equivalent points in its interior. The Poincarè algorithm (see [4], [6]) gives us for each edge segment class one cycle transformation of the form  $c = g_r \dots g_2 g_1$ , where  $g_i, i = 1, 2, \dots r$  are face pairing identifications. Since each of these cycle transformations is a rotation of order v, the cycle relations are of the form  $(g_r \dots g_2 g_1)^v = 1$ . The Poincarè theorem guarantees that these cycle relations, together with relations,  $g_i^2 = 1$  to the occasional involutive generators  $g_i = g_i^{-1}$ , form a complete set of defining relations for *G*.

**Remark 1.** As an intuitive picture (e.g. from [6]) we shall use in the following, the space tiling  $P^G$  by fundamental domains (polyhedra) represents the space group G itself after having chosen an arbitrary identity domain P. For any generator  $g \in G$  domain P will be adjacent with the g-image domain  $P^g$  along face  $f_g$  denoted simply either by g on the side (half space) of  $P^g$ , or by  $g^{-1}$  on the side of P (we prefer the first). Involutive generators  $g = g^{-1}$  we write just on the corresponding face. Thus for any  $\gamma \in G$ , domain  $P^{\gamma}$  will be adjacent with domain  $P^{g\gamma}$  along an image face  $f_g^{\gamma}$ , signed simply either by g on side of  $P^{g\gamma}$ , or by  $g^{-1}$  on the side of  $P^{\gamma}$ . So domain  $P^{\gamma}$  is obtained by  $\gamma = g_r \dots g_2 g_1$  from P through the path in domains P,  $P^{g_1}, P^{g_2 g_1}, \dots, P^{\gamma}$  entered finely through  $g_r$ . Imagine also the group graph by coloured directed edges with so many colours as many generators we have. From g and  $g^{-1}$  we use only g, as to direction of the corresponding graph edge.

**Classification of all non-fundamental trunc-simplex tilings** means the following finite process (7 steps as a natural sketchy generalization of our earlier works by *D*–symbol method or barycentric subsimplex orbits, see e.g. [10] and [11], you can skip this difficult part at the first read).

1. We consider a 3-simplex *T* (as tetrahedron) with its barycentric subdivision into 24 subsimplices. Any subsimplex has a 3-centre (solid centre, first formally, later also metrically), a 2-centre in a simplex face, a 1-centre in an edge of the previous face, a 0-centre as a vertex of the previous edge, as usual. The simplex has first the piece-wise linear (PL) topology. Any subsimplex has a 3-face. 2-, 1-, 0- faces, opposite to the corresponding centres. Then 2-, 1-, 0- (involutive (involutory)) adjacency operations  $\sigma^i$  can be introduced. Then occasional symmetry operations come acting on points, segments ..., preserving the subsimplices and incidences with adjacency operations. E.g. we can write by convention

$$\sigma^i C^g := (\sigma^i C)^g = \sigma^i (C^g), \tag{1.1}$$

that means for any subsimplex *C*: the adjacency operation  $\sigma^i$  ( $i \in \{0, 1, 2\}$  first) and mapping *g* commute. So,  $\sigma^i$ 's act on left, symmetries act on right.

- 2. The  $\sigma^i$  operation, as above ( $i \in \{0, 1, 2\}$  at the first considerations), can be considered as a local reflection in the *i*-side face of any subsimplex C. Then any symmetry mapping q (if it exists) can be given by a starting subsimplex and its image subsimplex and by simplex-wise extension, so that (1.1)holds further step-by-step, as requirement. The axiomatically defined symmetric D-matrix function  $(m^{ij} = m^{ji})$  and the consequences for *D*-symbol "theory" (e.g. in [10], [11]) can guarantee the extension of g first to the first simplex T, then – introducing  $\sigma^3$  operation for the tile adjacency in the next step – onto the whole tiling  $\mathcal{T}$ .
- 3. The face pairing generator symmetry mappings of *T*, i.e.  $g_1, g_2, \ldots$  (may be involutive one as well), introduce also  $\sigma^3$  adjacency operations and generate a simplex tiling  $\mathcal{T} = T^G$  under a symmetry group G, finitely generated by the previous  $g_1, g_2, \ldots$  The D-matrix function has free parameters yet for  $m^{23} = m^{32}$  entries, so we shall have infinite series of tilings for a given D-diagram with certain free rotational orders. The D-matrix function  ${}_{k}D \mapsto m^{ij}({}_{k}D)$  of the D-symbol says, by requirement, that the *ij*-edge at the meeting of  $\sigma^i$ ,  $\sigma^j$  side faces of any subsimplex  ${}_kC$  in the *D*-set of the *D*-symbol, schould be surrounded by  $2m^{ij}(_kD)$  subsimplices in the tiling  $\mathcal{T} = T^G$ . (Note, e.g. for the 3–simplex tiling  $\mathcal{T}$  that  $m^{01} = m^{12} = 3$  constant, and  $m^{02} = m^{03} = m^{13} = 2$  constant, because of the barycentric subdivision of  $\mathcal{T}$ .)
- 4. Around the half-edge G-equivalence classes we get subsimplices providing defining relations for G(besides the involution relations), just in the sense of Poincaré. However, it is our main consideration, we can take the  $D^0$  subsymbols, obtained by leaving  $\sigma^0$  operations from the D-diagram (and the D-matrix entries), the components describing the fundamental domains of the vertex stabilizer subgroups for the vertex equivalence classes induced, as a cutting 2-surface. The curvature formula ([10], [11]) says whether a vertex stabilizer (depending on the free parameters) is either a finite spherical group (of positive curvature), or Euclidean group (of zero curvature), or hyperbolic one (of negative curvature).
- 5. In this paper we concentrate on the last cases, where the parameters involve only outer vertices of hyperbolic space  $\mathbb{H}^3$ . Than we cut (truncate) this vertex(ices) orthogonally by polar plane(s) in  $\mathbb{H}^3$ . The vertex stabilizer group acts then, preserving the two half spaces of the polar plane(s). Introducing half-space changing isometries to this  $\mathbb{H}^2$ -stabilizer, preserving the original simplex tiling  $\mathcal{T}$ , so we get a space group with compact fundamental domain. The procedure induces face pairing at truncating vertex polars as well.
- 6. This new face pairings can be formulated by a new adjacency operation  $\sigma^{0*}$  commuting with other  $\sigma^{j'}$ s ( $j \in \{1, 2, 3\}$ ) but not with  $\sigma^0$ . Computer can also help, but we have made careful case-by-case discussions, first for later experiences.
- 7. A D-symbol isomorphism leads to equivariant group extension. Morphism does it onto a smaller D-diagram, or opposite procedures, lead to symmetry breaking members in the same family, where the free parameters need extra attention (see [10] and [11]).

#### 1.3. Results

In Section 2 we illustrate our method with some cases from family F1, then the leading representative group series will be detailed from each family (Sections 2-5), sometimes with characteristic examples, also accompanied with figures. Finally we have proved the following summary

**Theorem 1.** For integer parts of the trunc-simplices of families F1 - F4 there are given the 73 group extensions all together. Namely,

F1: Representing series  ${}^{*233}\Gamma(u) = {}^{\bar{4}3m}_{24}\Gamma(u)$  have 14 extensions, F2: Representing series  ${}^{2*2}\Gamma_3(u, 2v) = {}^{\bar{4}2m}_{8}\Gamma(u, 2v)$  have 21 extensions,

F3: Representing series  ${}^{*33}\Gamma(2u, v) = {}^{3m}_{6}\Gamma(2u, v)$  have 17 extensions, F4: Representing series  ${}^{*22}\Gamma_{1}(u, 2v, w) = {}^{mm2}_{4}\Gamma_{1}(u, 2v, w)$  have 21 extensions.

There are given all the space group series in figures and by group presentations, namely, the face pairing generators and defining relations to the edge-segment equivalence classes in Sections 2-5. These need standard but tedious work, of course. The analogous [10] other non-fundamental series for families F5 - F12 will be published later. Then we get a huge list of hyperbolic space groups in infinite series, depending on some parameters.

The truncation plane of a trunc-simplex will be a natural base plane for a hypersphere (constant distance surface from the base plane in both sides), whose ball (the inner part of the hypersphere) with its images under the symmetry group of the corresponding trunc-simplex tiling, can fil (pack) the whole space  $\mathbb{H}^3$  without common interior point, or can cover  $\mathbb{H}^3$  without gap, respectively.

As a by-product, we have determined the optimal, densest packing for each family from F1-F4, moreover, the optimal thinnest (loosest) covering for families F1-F2. Now this is straightforward, because the Coxeter-Schläfli (extended) reflection groups  $\{\overline{u}, \overline{v}, \overline{w}\}$  occur for these families. See tables to Section 6. However, the coverings for families F3-F4 deserve a separate publication with worser densities than the optimal covering at F2.

**Theorem 2.** Among the hyperball packings and coverings to the truncation base planes of hyperballs in the truncsimplex tilings in our families F1-F4 there is a densest packing in F1 to group  $^{*233}\Gamma$  (u = 7) (as complete Coxeter group {3,3,7}) with density  $\approx 0.82251$  (in Fig. 1); and there is a loosest covering in F2 to group  $^{2*2}\Gamma$  (u = 7, 2v = 6) (extended {7,3,7}) with density  $\approx 1.26829$  (in Fig. 8 and Fig. 25).

We conjecture (see also investigations [21–26] of the third author in References) that we have found among all arrangements of congruent hyperballs in  $\mathbb{H}^3$  the densest hyperball packing with density  $\approx 0.82251$  above (related to Dirichlet-Voronoi cell subdivision) to the regular trunc-simplex group {3,3,7}; furthermore the loosest hyperball covering with density  $\approx 1.26829$  above, to the extended trunc-orthoscheme group {7,3,7}. This last result is published first time here (surprising a little bit). This covering density is less than that in the conjecture of L. Fejes Tóth [1] for balls and horoballs.

## 2. Family F1

The maximal series of family F1 is characterized by the simplex tiling whose group series is  ${}^{43m}_{24}\Gamma(u)$  with previous (crystallographic) notation from [10], and now with the new (orbifold) notation  ${}^{*233}\Gamma(u)$  (see [3]). It has the trivial extension by the reflection  $a_0 = \bar{m}_0$  (Fig. 1: F1-1)

$${}^{*233}\Gamma(u) = (m_0, m_1, m_2, m_3 - 1 = m_0^2 = m_1^2 = m_2^2 = m_3^2 = (m_0 m_1)^3 = (m_0 m_2)^2 =$$
$$= (m_0 m_3)^2 = (m_1 m_2)^3 = (m_1 m_3)^2 = (m_2 m_3)^u; \ 6 < u \in \mathbb{N})$$
with extenion  $(a_0 - 1 = a_0^2 \stackrel{1}{=} (a_0 m_1)^2 \stackrel{2}{=} (a_0 m_2)^2 \stackrel{3}{=} (a_0 m_3)^2).$ 

# Other non-fundamental cases in F1:

•  ${}^{233}\Gamma_1(2u) = (r, r_2, m_3 - 1 = r^3 = r_2^2 = m_3^2 = m_3 r m_3 r^{-1} = (rr_2)^3 = (m_3 r_2 m_3 r_2)^u$ ,  $3 < u \in \mathbb{N}) = {}^{23}_{12}\Gamma_1(2u)$ , with 2 extensions (1)  $(\bar{m}_0, \bar{m}_1 - 1 = \bar{m}_0^2 = \bar{m}_1^2 \stackrel{1}{=} \bar{m}_0 r^{-1} \bar{m}_1 r \stackrel{2}{=} \bar{m}_0 r_2 \bar{m}_1 r_2 \stackrel{3}{=} (\bar{m}_0 m_3)^2 \stackrel{4}{=} (\bar{m}_1 m_3)^2)$ (2)  $(\bar{s} - 1 \stackrel{1}{=} (\bar{s}r)^2 \stackrel{2}{=} (\bar{s}r_2)^2 \stackrel{3}{=} m_3 \bar{s} m_3 \bar{s}^{-1})$ .

The pictures are given in Fig. 1: F1-2 with  $\Gamma^0(A_0, A_1) = 3 * u$  as  $\mathbb{H}^2$  group with half-turn extension.

•  ${}^{233}\Gamma_2(u) = (r, r_2, r_3 - 1 = r^3 = r_2^2 = r_3^2 = (rr_2)^3 = (rr_3)^2 = (r_2r_3)^u, 6 < u \in \mathbb{N}) = {}^{23}_{12}\Gamma_2(u)$ , with 2 extensions (1)  $(\bar{m}_0, \bar{m}_1 - 1 = \bar{m}_0^2 = \bar{m}_1^2 \stackrel{1}{=} \bar{m}_0 r^{-1} \bar{m}_1 r \stackrel{2}{=} \bar{m}_0 r_2 \bar{m}_1 r_2 \stackrel{3}{=} \bar{m}_0 r_3 \bar{m}_1 r_3)$ (2)  $(\bar{s} - 1 \stackrel{1}{=} (\bar{s}r)^2 \stackrel{2}{=} (\bar{s}r_2)^2 \stackrel{3}{=} (\bar{s}r_3)^2)$ .

In Fig. 2: F1-3 the pictures are given with  $\Gamma^0 = 23u$  as  $\mathbb{H}^2$  group with half-turn extension.

• 
$${}^{2\times}\Gamma_3(3u) = (z, r_1, r_3 - 1 = r_1^2 = r_3^2 = zzr_1 = (r_1r_3zr_3z^{-1}r_3)^u, 2 < u \in \mathbb{N}) = \frac{4}{4}\Gamma_3(3u)$$
, with 2 extensions

(1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 \stackrel{1}{=} \bar{m}_0 r_1 \bar{m}_2 r_1 \stackrel{2}{=} \bar{m}_0 z^{-1} \bar{m}_1 z \stackrel{3}{=} \bar{m}_0 r_3 \bar{m}_1 r_3 \stackrel{4}{=} \bar{m}_1 z^{-1} \bar{m}_2 z \stackrel{5}{=} (\bar{m}_2 r_3)^2)$ (2)  $(\bar{g}, \bar{h}_2 - 1 = \bar{h}_2^2 \stackrel{1}{=} \bar{g}z^{-1}\bar{h}_2r_1 \stackrel{2}{=} (\bar{g}z)^2 \stackrel{3}{=} (\bar{g}r_3)^2 \stackrel{5}{=} (\bar{h}_2r_3)^2$ . The pictures are given in Fig. 2: F1-4 with  $\Gamma^0 = 2u \times$  as  $\mathbb{H}^2$  group with half-turn extension.

•  ${}^{33}\Gamma_5(4u) = (r, r_2, r_3 - 1 = r^3 = r_2^2 = r_3^2 = (rr_2)^2 = (r_3 r r_3 r^{-1} r_3 r_2)^u, \frac{3}{2} < u \in \mathbb{N}) = {}^{3}_{3}\Gamma_5(4u)$ , with 2 extensions (1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 \stackrel{1}{=} \bar{m}_0 r^{-1} \bar{m}_1 r \stackrel{2}{=} \bar{m}_0 r_2 \bar{m}_1 r_2 \stackrel{3}{=}$  $\bar{m}_0 r_3 \bar{m}_2 r_3 \stackrel{4}{=} (\bar{m}_1 r_3)^2 \stackrel{5}{=} \bar{m}_2 r \bar{m}_2 r^{-1})$ (2)  $(\bar{h}_0, \bar{h}_1, \bar{h}_2 - 1 = \bar{h}_0^2 = \bar{h}_1^2 = \bar{h}_2^2 \stackrel{1}{=} \bar{h}_0 r_2 \bar{h}_1 r \stackrel{3}{=} \bar{h}_0 r_3 \bar{h}_2 r_3 \stackrel{4}{=} (\bar{h}_1 r_3)^2 \stackrel{5}{=} (\bar{h}_2 r)^2).$ The pictures with  $\Gamma^0 = 223$  as  $\mathbb{H}^2$  group with two half-turn extensions are given in Fig. 3: F1-5.

•  ${}^{33}\Gamma_6(8u) = (r, m_2, r_3 - 1 = r^3 = m_2^2 = r_3^2 = m_2 r m_2 r^{-1} = (m_2 r_3 r r_3 r^{-1} r_3)^{2u}, \frac{3}{4} < u \in \mathbb{N}) = {}^{3}_{3}\Gamma_6(8u)$ , with trivial extension

(1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 \stackrel{1}{=} \bar{m}_0 r^{-1} \bar{m}_1 r \stackrel{2}{=} (\bar{m}_0 m_2)^2 \stackrel{3}{=} \bar{m}_0 r_3 \bar{m}_2 r_3 \stackrel{4}{=} (\bar{m}_1 m_2)^2 \stackrel{5}{=} (\bar{m}_1 r_3)^2 \stackrel{6}{=} \bar{m}_2 r \bar{m}_2 r^{-1}).$ See Fig. 3: F1-6 for the pictures with  $\Gamma^0 = 23 * u$ .

•  ${}^{22}\Gamma_{11}(6u) = (z, h - 1 = h^2 = (hzzhzhz^{-1}hz^{-1}z^{-1})^u, \ 1 < u \in \mathbb{N}) = {}^2_2\Gamma_{11}(6u),$ 

with 4 extensions, each in two variants, given in Fig. 4 - 6 to F1-7. These so-called Gieseking orbifolds are derived from the famous non-orientable manifold with an infinite vertex class, originally if u = 1 (cusp, but 1 < u now). The half-turn extension h allows various half simplices and fundamental domains, e.g. for the trivial extension as follows (Fig. 4: F1-7<sub>1</sub>).

(1) 
$$(\bar{m}_0, \bar{m}_1, \bar{m}_2, \bar{m}_3 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 = \bar{m}_3^2 = \bar{m}_0 h \bar{m}_1 h \stackrel{?}{=} \bar{m}_0 z^{-1} \bar{m}_1 z \stackrel{?}{=} \bar{m}_1 z^{-1} \bar{m}_2 z \stackrel{4}{=} \bar{m}_2 h \bar{m}_3 h \stackrel{?}{=} \bar{m}_3 z \bar{m}_3 z^{-1}).$$
  
Fig. 5: F1-7<sub>2</sub> to <sup>22</sup> $\Gamma_{11}(6u)$  show the first version of extensions (2)-(4) symbolically.

Fig. 6: F1-7<sub>3</sub> show the simpler version of extensions (2)-(4) by (additional) presentations:

(2) 
$$(\bar{h}_0, \bar{h}_2 - 1 = \bar{h}_0^2 = \bar{h}_2^2 \stackrel{2}{=} (\bar{h}_0 z^{-1} h)^2 \stackrel{3}{=} \bar{h}_0 h z^{-1} \bar{h}_2 z h \stackrel{5}{=} (\bar{h}_2 h z h)^2)$$

(3)  $(\bar{q} - 1 \stackrel{2}{=} \bar{q}hzh\bar{q}^{-1}z^{-1}h \stackrel{3}{=} (\bar{q}zh)^2),$ (4)  $(\hat{s} - 1 \stackrel{2}{=} \hat{s}hzh\hat{s}^{-1}hz \stackrel{3}{=} (\hat{s}zh)^2)$ .

### 2.1. The 4 extensions to Gieseking orbifolds as fundamental trunc-simplex tilings

Our Fig. 7: F1-7+ to  $\Gamma_{62}(6u)$  shows this famous simplex face pairing, as repetition from [20], with orientation reversing transforms

$$\mathbf{z}_1: z_1^{-1} \to z_1$$
 fixing vertex  $A_3$  and  $\mathbf{z}_2: z_2^{-1} \to z_2$  fixing  $A_0$ 

All vertices lie at the infinity (absolute) of  $\mathbb{H}^3$ , iff the "rotation order" parameter at the edge equivalence class is u = 1. At the polar plane of  $A_0$  (say, e.g.) we look the fundamental domain  $F^0$  of the vertex class stabilizer subgroup  $\Gamma^0 = uu \times x$  that is the Euclidean Klein bottle group **4.pg** (with two cross caps  $\times$ , or projective planes), iff u = 1. Moreover, we get interesting "surgery effects" that lead e.g. to the famous Fomenko-Matveev-Weeks hyperbolic space form of minimal volume [27].

Otherwise, iff  $1 < u \in \mathbb{N}$  we get  $\mathbb{H}^2$  plane group (area of  $F^0$  is  $\frac{4\pi}{u}$  less than  $4\pi$  [3]). Then the generators, reversing the half-spaces of the polar plane  $a_0$ , but preserving the original simplex tiling T, lead to 4 group extensions with compact fundamental domains as trunc-simplices (or octahedra  $O_{62}^1, \ldots, O_{62}^4$ ).

Presentation of the original simplex tiling group is

$$\Gamma_{62}(6u) = (z_1, z_2 - 1) = (z_1 z_2 z_2 z_1^{-1} z_2^{-1} z_1)^u; \ 1 < u \in \mathbb{N}).$$

(1) Then comes the trivial extension with plane reflections:

$$(\overline{m}_0, \overline{m}_1, \overline{m}_2, \overline{m}_3 - 1 = \overline{m}_0^2 = \overline{m}_1^2 = \overline{m}_2^2 = \overline{m}_3^2 \stackrel{1}{=} \overline{m}_0 z_1^{-1} \overline{m}_2 z_1 \stackrel{2}{=} \overline{m}_0 z_2 \overline{m}_0 z_2^{-1} \stackrel{3}{=} \stackrel{3}{=} \overline{m}_1 z_1 \overline{m}_2 z_1^{-1} \stackrel{4}{=} \overline{m}_1 z_2 \overline{m}_2 z_2^{-1} \stackrel{5}{=} \overline{m}_1 z_2^{-1} \overline{m}_3 z_2 \stackrel{6}{=} \overline{m}_3 z_1 \overline{m}_3 z_1^{-1}).$$

(2) The half turns extension, changing the half spaces at  $F^0$  (inducing point reflections at edges 2,6, moreover, a 2–knot *u*–segment combination):

$$(\overline{h}_0, \overline{h}_1, \overline{h}_2, \overline{h}_3 - 1 = \overline{h}_0^2 = \overline{h}_1^2 = \overline{h}_2^2 = \overline{h}_3^2 \stackrel{1}{=} \overline{h}_0 z_1^{-1} \overline{h}_2 z_1 \stackrel{2}{=} (\overline{h}_0 z_1)^2 \stackrel{3}{=} \frac{3}{\overline{h}_1 z_1 \overline{h}_2 z_2^{-1}} \stackrel{5}{=} \overline{h}_1 z_2^{-1} \overline{h}_3 z_2 \stackrel{6}{=} (\overline{h}_3 z_1)^2.$$

(3) The most interesting half-screw extension, changing the half-spaces at  $F^0$  yields a "non-orientable *u*-knot" (by glide reflections  $\overline{g}_1, \overline{g}_2$ ):

$$(\overline{g}_1, \overline{g}_2 - 1 \stackrel{1}{=} \overline{g}_1 z_2 \overline{g}_2^{-1} z_1 \stackrel{2}{=} \overline{g}_1 z_1 g_1^{-1} z_2 \stackrel{3}{=} \overline{g}_2 z_1 \overline{g}_2 z_2)$$

(4) The 4th extension is a point reflection at edges 3, 4 of  $F^0$  (inducing screw motions  $\overline{s}_1, \overline{s}_2$  and u-knot):

$$(\bar{s}_1, \bar{s}_2 - 1 \stackrel{1}{=} \bar{s}_1 z_2 \bar{s}_2^{-1} z_1 \stackrel{2}{=} \bar{s}_1 z_1 s_1^{-1} z_2^{-1} \stackrel{3}{=} (\bar{s}_2 z_1)^2 \stackrel{4}{=} (\bar{s}_2 z_2)^2).$$

We think that these phenomena are remarkable!

# 3. Family 2

In family F2 the maximal series is characterized by the simplex tiling whose group series is  $2^{*2}\Gamma(u, 2v)$  (or  $\frac{42m}{8}\Gamma(u, 2v)$  by notations in [10]). Remember that in case u = 2v we order this tiling (and similarly the later ones) to family F1 with parameter  $\overline{u} = u = 2v$ , since so many simplices (i.e.  $2\overline{u}$  barycentric subsimplices) surround every simplex edge.

$${}^{2*2}\Gamma(u,2v) = (m_0,m_2,m_3,r_1-1=m_0^2=m_2^2=m_3^2=r_1^2=(m_0m_2)^2=(m_0m_3)^2=$$
$$= m_0r_1m_2r_1 = (m_2m_3)^u = (m_3r_1m_3r_1)^v, \ u \neq 2v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < \frac{1}{2})$$

with the trivial extension

 $(\bar{m}_0, \bar{m}_2 - 1 = \bar{m}_0^2 = \bar{m}_2^2 \stackrel{1}{=} \bar{m}_0 r_1 \bar{m}_2 r_1 \stackrel{2}{=} (\bar{m}_0 m_2)^2 \stackrel{3}{=} (\bar{m}_0 m_3)^2 \stackrel{4}{=} (\bar{m}_2 m_0)^2 \stackrel{5}{=} (\bar{m}_2 m_3)^2).$ 

Pictures are given in Fig. 8: F2-1 with  $\Gamma^0(A_0, A_2) = *2uv$  as  $\mathbb{H}^2$  group.

### Other non-fundamental cases in F2

• \*<sup>22</sup> $\Gamma_2(2u, 2v) = (m_0, m_2, z - 1 = m_0^2 = m_2^2 = (m_0m_2)^2 = m_2 z m_0 z^{-1} =$  $(m_0 z m_2 z^{-1})^u = (zz)^v, \ u \neq v \in \mathbb{N}, \ \frac{1}{u} + \frac{2}{v} < 1) = \frac{mm_2^2}{4} \Gamma_2(2u, 2v)$ with 2 extensions in Fig. 8: F2-2 (1)  $(\bar{m}_0, \bar{m}_2 - 1 = \bar{m}_0^2 = \bar{m}_2^2 \stackrel{1}{=} \bar{m}_0 z \bar{m}_2 z^{-1} \stackrel{2}{=} (\bar{m}_0 m_2)^2 \stackrel{3}{=} \bar{m}_0 z^{-1} \bar{m}_2 z \stackrel{4}{=} (\bar{m}_2 m_0)^2)$ (2)  $(\bar{g} - 1 \stackrel{1}{=} (\bar{g}z^{-1})^2 \stackrel{2}{=} m_2 \bar{g} m_0 \bar{g}^{-1} \stackrel{3}{=} (\bar{g}z)^2)$ .  $\Gamma^0 = v * u$  is a  $\mathbb{H}^2$  group with half-turn extension yielding glide reflection  $\bar{g}$ .

•  ${}^{2\times}\Gamma_1(2u,2v) = (r_2,m_3,z-1) = r_2^2 = m_3^2 = zzr_2 = (m_3r_2m_3r_2)^u = (m_3zm_3z^{-1})^v, \ u \neq v \in \mathbb{N}, \ \frac{1}{u} + \frac{2}{v} < 1) = r_2^2$  $\frac{1}{4}\Gamma_1(2u, 2v)$  with 2 extensions

(1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 \stackrel{1}{=} \bar{m}_0 z^{-1} \bar{m}_2 z \stackrel{2}{=} \bar{m}_0 r_2 \bar{m}_1 r_2 \stackrel{3}{=} (\bar{m}_0 m_3)^2 \stackrel{4}{=} \bar{m}_1 z \bar{m}_2 z^{-1} \stackrel{5}{=} (\bar{m}_1 m_3)^2 \stackrel{6}{=} (\bar{m}_2 m_3)^2)$ (2)  $(\bar{h}_2, \bar{s} - 1 = \bar{h}_2^2 \stackrel{1}{=} \bar{s}z\bar{h}_2 z \stackrel{2}{=} (\bar{s}r_2)^2 \stackrel{3}{=} m_3 \bar{s}m_3 \bar{s}^{-1} \stackrel{6}{=} (m_3 \bar{h}_2)^2).$ Pictures are given in Fig. 9: F2-3 with  $\Gamma^0 = *uvv$  as  $\mathbb{H}^2$  group with half-turn extension.



Figure 1: F1-1 to  ${}^{*233}\Gamma(u)$  and F1-2 to  ${}^{233}\Gamma_1(2u)$ 



Figure 2: F1-3 to  $^{233}\Gamma_2(u)$  and F1-4 to  $^{2\times}\Gamma_3(3u)$ 



Figure 3: F1-5 to  ${}^{33}\Gamma_5(4u)$  and F1-6 to  ${}^{33}\Gamma_6(8u)$ 





 $\begin{array}{l} & \stackrel{-2}{}_{23} \\ h \text{ axis through } A_{01}A_{23} \text{ and simplex centre } A_{0123} \\ z \xrightarrow{h} z^{h} = \bar{z} = hzh \\ A_{l}^{h} = A_{0}, A_{3}^{h} = A_{2}, \text{ the side faces of } h \text{ can be varied} \\ This is a modified version of } \Gamma_{62}(6u) \end{array}$ 

Figure 4: F1-7<sub>1</sub> to  ${}^{22}\Gamma_{11}(6u)$  as Gieseking orbifolds



Figure 5: F1-7<sub>2</sub> to  ${}^{22}\Gamma_{11}(6u)$ , first version of ext. (2)-(4)



Figure 6: F1-7<sub>3</sub> to  ${}^{22}\Gamma_{11}(6u)$ , second simpler version of ext. (2)-(4)



Figure 7: F1-7+ to  $\Gamma_{62}(6u)$ , the 4 extensions to Gieseking fundamental orbifolds

•  ${}^{2\times}\Gamma_2(u,2v) = (r_2,r_3,z-1=r_2^2=r_3^2=zzr_2=(r_2r_3)^u = (zr_3zr_3)^v, u \neq 2v \in \mathbb{N}, \frac{1}{u} + \frac{1}{v} < \frac{1}{2}) = \frac{4}{4}\Gamma_2(u,2v)$  with 2 extensions (1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 \stackrel{1}{=} \bar{m}_0 z^{-1} \bar{m}_2 z \stackrel{2}{=} \bar{m}_0 r_2 \bar{m}_1 r_2 \stackrel{3}{=} \bar{m}_0 r_3 \bar{m}_1 r_3 \stackrel{4}{=} \bar{m}_1 z \bar{m}_2 z^{-1} \stackrel{5}{=} (\bar{m}_2 r_3)^2)$ (2)  $(\bar{h}_2, \bar{s} - 1 = \bar{h}_2^2 \stackrel{1}{=} \bar{s}z\bar{h}_2z \stackrel{2}{=} (\bar{s}r_2)^2 \stackrel{3}{=} (\bar{s}r_3)^2 \stackrel{5}{=} (\bar{h}_2r_3)^2).$ Pictures are given in Fig. 9: F2-4 with  $\Gamma^0 = 2uv$  as  $\mathbb{H}^2$  group with half-turn extension.

•  $^{222}\Gamma_2(u, 2v) = (r_0, r_1r_2, r_3 - 1 = r_0^2 = r_1^2 = r_2^2 = r_3^2 = r_0r_1r_2 = (r_2r_3)^u = (r_0r_3r_1r_3)^v, \ u \neq 2v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < \frac{1}{2}) = (r_0r_3r_1r_3)^v$  ${}^{222}_{4}\Gamma_2(u, 2v)$  with 2 extensions (1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 \stackrel{1}{=} \bar{m}_0 r_1 \bar{m}_2 r_1 \stackrel{2}{=} \bar{m}_0 r_2 \bar{m}_1 r_2 \stackrel{3}{=}$  $\bar{m}_0 r_3 \bar{m}_1 r_3 \stackrel{4}{=} \bar{m}_1 r_0 \bar{m}_2 r_1 \stackrel{5}{=} (\bar{m}_2 r_3)^2)$ (2)  $(\bar{h}_2, \bar{s} - 1 = \bar{h}_2^2 \stackrel{1}{=} \bar{s}r_0\bar{h}_2r_1 \stackrel{2}{=} (\bar{s}r_2)^2 \stackrel{3}{=} (\bar{s}r_3)^2 \stackrel{5}{=} (\bar{h}_2r_3)^2).$ 

See Fig. 10: F2-5 with  $\Gamma^0 = 2uv$  as  $\mathbb{H}^2$  group with half-turn extension.

# 3.1. Truncated half-simplices from F2

•  ${}^{22}\Gamma_7(2u, 4v) = (r_2, r_3, h - 1 = r_2^2 = r_3^2 = h^2 = (r_2r_3)^u = (r_2hr_3hr_3hr_2h)^v$ ,  $u \neq 2v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < 1 = \frac{2}{2} \Gamma_7(2u, 4v)$  with 4 extensions (1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2, \bar{m}_3 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 = \bar{m}_3^2 \stackrel{1}{=} \bar{m}_0 h \bar{m}_3 h \stackrel{2}{=} \bar{m}_0 r_2 \bar{m}_1 r_2 \stackrel{3}{=} \bar{m}_0 r_3 \bar{m}_1 r_3 \stackrel{4}{=} \bar{m}_1 h \bar{m}_2 h \stackrel{5}{=} (\bar{m}_2 r_3)^2 \stackrel{6}{=} (\bar{m}_3 r_2)^2)$ (2)  $(\bar{q}_1, \bar{q}_2 = 1 \stackrel{1}{=} \bar{q}_1 h \bar{q}_2 h \stackrel{2}{=} \bar{q}_1 r_3 \bar{q}_1 r_2 \stackrel{5}{=} r_3 \bar{q}_2 r_2 \bar{q}_2^{-1})$ (3)  $(\bar{s}_1, \bar{s}_2 = 1 \stackrel{1}{=} \bar{s}_1 h \bar{s}_2 h \stackrel{2}{=} (\bar{s}_1 r_2)^2 \stackrel{3}{=} (\bar{s}_1 r_3)^2 \stackrel{5}{=} r_3 \bar{s}_2 r_2 \bar{s}_2^{-1})$ (4)  $(\bar{h}_0, \bar{h}_1, \bar{h}_2, \bar{h}_3 - 1 = \bar{h}_0^2 = \bar{h}_1^2 = \bar{h}_2^2 = \bar{h}_3^2 \stackrel{1}{=} \bar{h}_0 h \bar{h}_3 h \stackrel{2}{=} \bar{h}_0 r_2 \bar{h}_1 r_3 \stackrel{4}{=} \bar{h}_1 h \bar{h}_2 h \stackrel{5}{=} (\bar{h}_2 r_3)^2 \stackrel{6}{=} (\bar{h}_3 r_2)^2).$ Look at Fig. 10: F2-6<sub>1</sub> for these presentations with  $\Gamma^0 = 22uv$  as  $\mathbb{H}^2$  group. At Fig. 11 : F2-6<sub>2</sub> you look

pictures for the simpler versions. E.g. for the third extension (by a half-turn about edge 3, so 2 as well, yielding screw motion  $\bar{s}$ ) we get in addition:

(3)  $(\bar{s} - 1 \stackrel{2}{=} (\bar{s}r_2)^2 \stackrel{3}{=} (\bar{s}r_3)^2 \stackrel{5}{=} \bar{s}hr_3h\bar{s}^{-1}hr_2h).$ 

•  ${}^{22}\Gamma_8(u, 4v) = (r, h-1 = h^2 = r^u = (hrhrhr^{-1}hr^{-1})^v, 2 < u \neq 4v \in \mathbb{N}, \frac{2}{u} + \frac{1}{v} < 1) = {}^2_2\Gamma_8(u, 4v)$  with 4 extensions (1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2, \bar{m}_3 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 = \bar{m}_3^2 \stackrel{1}{=} \bar{m}_0 h \bar{m}_3 h \stackrel{2}{=} \bar{m}_0 r \bar{m}_0 r^{-1} \stackrel{3}{=} \bar{m}_1 h \bar{m}_2 h \stackrel{4}{=} \bar{m}_1 r \bar{m}_1 r^{-1} \stackrel{5}{=} \bar{m}_2 r^{-1} \bar{m}_3 r)$ (2)  $(\bar{s}_1, \bar{s}_2 = 1 \stackrel{1}{=} \bar{s}_1 h \bar{s}_2 h \stackrel{2}{=} \bar{s}_1 r \bar{s}_1^{-1} r^{-1} \stackrel{5}{=} (r \bar{s}_2)^2)$ (3)  $(\bar{h}_0, \bar{h}_1, \bar{h}_2, \bar{h}_3 - 1 = \bar{h}_0^2 = \bar{h}_1^2 = \bar{h}_2^2 = \bar{h}_3^2 \stackrel{1}{=} \bar{h}_0 h \bar{h}_3 h \stackrel{2}{=} (\bar{h}_0 r)^2 \stackrel{3}{=} \bar{h}_1 h \bar{h}_2 h \stackrel{4}{=} (\bar{h}_1 r)^2 \stackrel{5}{=} \bar{h}_2 r^{-1} \bar{h}_3 r)$ (4)  $(\bar{g}_1, \bar{g}_2 = 1 \stackrel{1}{=} \bar{g}_1 h \bar{g}_2 h \stackrel{2}{=} \bar{g}_1 r \bar{g}_1^{-1} r \stackrel{5}{=} (r \bar{g}_2)^2).$ Look at Fig. 12: F2-7 with  $\Gamma^0 = uuv$  as  $\mathbb{H}^2$  group. Simpler presentation is not pictured.

•  ${}^{22}\Gamma_9(2u, 2v) = (s, h - 1 = h^2 = (shs^{-1}h)^u = (ssh)^v = 1, 1 < u \in \mathbb{N}, 1 < v \in \mathbb{N}, \frac{1}{u} + \frac{2}{v} < 1) = {}^2_2 \Gamma_9(2u, 2v)$  with 2 extensions (1)  $(\bar{m}_0, \bar{m}_1, \bar{m}_2, \bar{m}_3 - 1 = \bar{m}_0^2 = \bar{m}_1^2 = \bar{m}_2^2 = \bar{m}_3^2 \stackrel{1}{=} \bar{m}_0 h \bar{m}_1 h \stackrel{2}{=} \bar{m}_0 s^{-1} \bar{m}_3 s \stackrel{3}{=} \bar{m}_1 s \bar{m}_2 s^{-1} \stackrel{4}{=} \bar{m}_1 s^{-1} \bar{m}_2 s \stackrel{5}{=} \bar{m}_2 h \bar{m}_3 h)$ 

(2)  $(\bar{g}_1, \bar{g}_2 = 1 \stackrel{1}{=} \bar{g}_1 h \bar{g}_2^{-1} h \stackrel{2}{=} \bar{g}_1 s \bar{g}_2 s \stackrel{3}{=} (\bar{g}_2 s)^2).$ 

Look at Fig. 13: F2-8<sub>1</sub> for these presentations with  $\Gamma^0 = uvv$  as  $\mathbb{H}^2$  group. For the simpler second version look at Fig. 14 F2-82. For instance the second extension occurs with point reflection at edge 3, yielding glide reflection  $\bar{g}$  as follows in addition:

(2) 
$$(\bar{g} - 1 \stackrel{2}{=} \bar{g}sh\bar{g}hs \stackrel{3}{=} (\bar{g}hs^{-1}h)^2).$$

1422

### 3.2. The 8 extensions for fundamental trunc-simplex tiling and group $\Gamma_{58}(2u, 4v)$

These top number extensions are pictured in Fig. 15 and 16 as to F2-8+. Maybe, the numbered edges at truncation are more convenient than in the original publication [19]. The stabilizer subgroup  $\Gamma^0 = uuvv$  of the vertex class are pictured in Fig. 15, where every vertex triangle has one *u*-point ( $u^+$  or  $u^-$ ) and  $v^+$ -, $v^-$ -point as vertices to the spatial rotation axes, respectively. Thus the 8 possible correspondences are upper bound, indeed, that can be realized.

$$\Gamma_{58}(2u, 4v) = (s_1, s_2 - 1) = (s_1 s_2^{-1})^u = (s_1 s_1 s_2 s_2)^v, \ u \neq 2v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < 1).$$

(1) The trivial extension is for  $O_{58}^1$  as octahedron:

$$(\overline{m}_{0}, \overline{m}_{1}, \overline{m}_{2}, \overline{m}_{3} - 1 = \overline{m}_{0}^{2} = \overline{m}_{1}^{2} = \overline{m}_{2}^{2} = \overline{m}_{3}^{2} \stackrel{1}{=} \overline{m}_{0} s_{2} \overline{m}_{2} s_{2}^{-1} \stackrel{2}{=} \overline{m}_{0} s_{1}^{-1} \overline{m}_{3} s_{1} \stackrel{3}{=} \stackrel{3}{=} \overline{m}_{0} s_{2}^{-1} \overline{m}_{3} s_{2} \stackrel{4}{=} \overline{m}_{1} s_{1} \overline{m}_{3} s_{1}^{-1} \stackrel{5}{=} \overline{m}_{1} s_{1}^{-1} \overline{m}_{2} s_{1} \stackrel{6}{=} \overline{m}_{1} s_{2}^{-1} \overline{m}_{2} s_{2}).$$

(2) The second extension will be derived from the two half-turns at  $u^+$  and  $u^-$  that fix the vertex triangles by  $\bar{h}_0, \ldots, \bar{h}_3$  as follows for  $O_{58}^2$  in addition:

$$(\bar{h}_0, \bar{h}_1, \bar{h}_2, \bar{h}_3 - 1 = \bar{h}_0^2 = \bar{h}_1^2 = \bar{h}_2^2 = \bar{h}_3^2 \stackrel{!}{=} \bar{h}_0 s_2 \bar{h}_2 s_2^{-1} \stackrel{!}{=} \bar{h}_0 s_1^{-1} \bar{h}_3 s_2 \stackrel{!}{=} \frac{4}{\bar{h}_1 s_1 \bar{h}_3 s_1^{-1}} \stackrel{!}{=} \bar{h}_1 s_1^{-1} \bar{h}_2 s_2).$$

(3) The 3rd extension will be by rotatory reflection at  $v^-$  (and  $v^+$  as well) rotating  $A_3^{s_1} \to A_2^{s_2^{-1}}$ ,  $A_0 \to A_1^{s_2^{-2}}$  and changing the half-spaces at  $F^0$ . This leads to glide reflection  $\overline{g}_1, \overline{g}_2$  and additional presentation:

$$(\overline{g}_1, \overline{g}_2 - 1 \stackrel{1}{=} \overline{g}_1 s_1 \overline{g}_2^{-1} s_2^{-1} \stackrel{2}{=} \overline{g}_1 s_2^{-1} g_2 s_1 \stackrel{3}{=} \overline{g}_1 s_1^{-1} \overline{g}_2 s_2).$$

(4) The 4th extension is derived by a half screw motions through D and  $\textcircled{D} A_2^{\overline{s_2}^{-1}} \rightarrow A_3^{\overline{s_1}}, A_0 \rightarrow A_1^{\overline{s_1}}$  changing the half-spaces at  $F^0$ . This leads to screw motions  $\overline{s_1}, \overline{s_2}$  and additional presentation:

$$(\bar{s}_1, \bar{s}_2 - 1 \stackrel{1}{=} \bar{s}_1 s_1 \bar{s}_2^{-1} s_2^{-1} \stackrel{2}{=} \bar{s}_1 \bar{s}_1^{-1} \bar{s}_2 s_1 \stackrel{3}{=} \bar{s}_1 \bar{s}_2^{-1} \bar{s}_2 s_2)^2)$$

Here we meet with fixed-point-free actions beyond the u- and v- rotations whose axes close. Thus we found a double (u, v) link as an interesting topological phenomenon. (The previous and next ones are also remarkable). E.g. the case u = 3, v = 2 leads to hyperbolic metric.

The case u = 2 = v leads to ideal vertices at the absolute (with  $\mathbb{E}^2$  metric to  $F^0$  domain, Fig. 15), i.e. we get  $\mathbb{H}^3$  realization with so-called cusp phenomenon. Imagine a fundamental tetrahedron in  $\mathbb{H}^3$  with two rectangles at opposite u = 2 rotational edges and  $\pi/4$  angles at the remaining four v = 2 rotational edges. We also cite from [10] the (u, 1) link cases with spherical  $\mathbb{S}^3$  realization. The "pure topological" (1, 1) link, as manifold with "lens realization", moreover, the 4–truncated lens as  $\mathbb{S}^3$  manifold are extremely remarkable, also to our context and Fig. 16.

(5) The 5th extension is obtained by point-reflection in (5) at edge 1 (or in (5) at 4):  $A_1^{s_2^{-2}} \rightarrow A_3^{s_2}, A_0 \rightarrow A_2^{s_2^{-1}}$ . These lead  $\widehat{g_1}, \widehat{g_2}$  pairings:

$$(\widehat{g}_1, \widehat{g}_2 - 1 \stackrel{1}{=} (\widehat{g}_1 s_2^{-1})^2 \stackrel{2}{=} \widehat{g}_1 s_1 \widehat{g}_2 s_1 \stackrel{3}{=} \widehat{g}_1 s_2 \widehat{g}_2 s_2) \stackrel{4}{=} (\widehat{g}_2 s_1^{-1})^2).$$

(6) The 6th extension is by half turn about edge 1 (also about 4)  $A_3^{s_1} \rightarrow A_1^{s_2^{-2}}$ ,  $A_0 \rightarrow A_2^{s_2^{-1}}$ , leading to screw motions  $\widehat{s_1}, \widehat{s_2}$ :

$$(\widehat{s}_1, \widehat{s}_2 - 1 \stackrel{1}{=} (\widehat{s}_1 s_2^{-1})^2 \stackrel{2}{=} \widehat{s}_1 s_2 \widehat{s}_2 s_1 \stackrel{3}{=} \widehat{s}_1 s_1 \widehat{s}_2 s_2) \stackrel{4}{=} (\widehat{s}_2 s_1^{-1})^2).$$

(7) The 7th extension will be  $A_0 \rightarrow A_3^*, A_1^* \rightarrow A_2^*$  by rotatory reflection at  $u^-$  and  $u^+$  (1  $\leftrightarrow$  4, 2  $\leftrightarrow$  3, 5  $\leftrightarrow$  6) yielding glide reflections  $\tilde{g}_1, \tilde{g}_2$ :

$$(\widetilde{g}_1, \widetilde{g}_2 - 1 \stackrel{1}{=} \widetilde{g}_1 s_1^{-1} \widetilde{g}_2 s_2^{-1} \stackrel{2}{=} \widetilde{g}_1 s_2 \widetilde{g}_1 s_1 \stackrel{5}{=} \widetilde{g}_2 s_1 \widetilde{g}_2 s_2).$$

(8) The 8th extension is again  $A_0 \rightarrow A_3^*, A_1^* \rightarrow A_2^*$ , but by half turns about edges 2, 3, 5, 6 with  $1 \leftrightarrow 4$  inducing screw motions  $\tilde{s_1}, \tilde{s_2}$ :

$$(\tilde{s}_1, \tilde{s}_2 - 1 \stackrel{1}{=} \tilde{s}_1 s_1^{-1} \tilde{s}_2 s_2^{-1} \stackrel{2}{=} (\tilde{s}_1 s_1)^2 \stackrel{3}{=} (\tilde{s}_1 s_2)^2 \stackrel{5}{=} (\tilde{s}_2 s_1)^2 \stackrel{6}{=} (\tilde{s}_2 s_2)^2).$$

# 4. Family F3

This family is also characterized by a Coxeter reflection group { $\bar{u}, \bar{v}, 3$ } with Coxeter Schläfli diagram in Fig. 17: F3-1 with doubly truncated simplex (trunc-orthoscheme). • \*<sup>33</sup> $\Gamma(2u, v) = (m_0, m_1, m_2, m_3 - 1 = m_0^2 = m_1^2 = m_2^2 = m_3^2 = (m_0 m_1)^3 =$ =  $(m_0 m_2)^2 = (m_0 m_3)^2 = (m_1 m_2)^2 = (m_1 m_3)^v = (m_2 m_3)^u$ ;  $2u \neq v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < \frac{1}{2}, \ \frac{1}{v} + \frac{1}{3} + \frac{1}{2} < 1$ ) =  ${}^{3m}_{6}\Gamma(2u, v)$ , with trivial extensions both for vertices  $A_0$  and  $A_2$ (1) for  $A_0$ :  $(\bar{m}_0 - 1 = \bar{m}_0^2 \stackrel{1}{=} (\bar{m}_0 m_1)^2 \stackrel{2}{=} (\bar{m}_0 m_2)^2 \stackrel{3}{=} (\bar{m}_0 m_3)^2$ ). (1) for  $A_2$ :  $(\bar{m}_2 - 1 = \bar{m}_2^2 \stackrel{4}{=} (\bar{m}_2 m_0)^2 \stackrel{5}{=} (\bar{m}_2 m_1)^2 \stackrel{6}{=} (\bar{m}_2 m_3)^2$ ). The pictures are given in Fig. 17: F3-1 with  $\Gamma^0(A_0) = *2uv$  and  $\Gamma^0(A_2) = *23v$  as  $\mathbb{H}^2$  groups.

•  ${}^{33}\Gamma_1(2u, 2v) = (r, m_2, m_3 - 1 = m_2^2 = m_3^2 = r^3 = (m_2m_3)^u = (m_3rm_3r^{-1})^v = m_2rm_2r^{-1}, \ u \neq v \in \mathbb{N}, \ \frac{2}{u} + \frac{1}{v} < 1, \ \frac{2}{3} + \frac{1}{v} < 1) = {}^{3}_{3}\Gamma_1(2u, 2v), \text{ with 2 extensions both for vertices } A_0, A_1 \text{ and } A_2$ (1) for  $A_0, A_1: (\bar{m}_0, \bar{m}_1 - 1 = \bar{m}_0^2 = \bar{m}_1^2 \stackrel{1}{=} \bar{m}_1r\bar{m}_0r^{-1} \stackrel{2}{=} (\bar{m}_0m_2)^2 \stackrel{3}{=} (\bar{m}_0m_3)^2 \stackrel{4}{=} (\bar{m}_1m_2)^2 \stackrel{5}{=} (\bar{m}_1m_3)^2)$ (1) for  $A_2: (\bar{m}_2 - 1 = \bar{m}_2^2 \stackrel{6}{=} \bar{m}_2r\bar{m}_2r^{-1} \stackrel{7}{=} (\bar{m}_2m_3)^2)$ (2) for  $A_0, A_1: (\bar{s} - 1 \stackrel{1}{=} (sr)^2 \stackrel{2}{=} \bar{s}m_2\bar{s}^{-1}m_2 \stackrel{3}{=} \bar{s}m_3\bar{s}^{-1}m_3)$ (2) for  $A_2: (\bar{h}_2 - 1 = \bar{h}_2^2 \stackrel{6}{=} (\bar{h}_2r)^2 \stackrel{7}{=} \bar{h}_2m_3\bar{h}_2m_3).$ The pictures are given in Fig. 17: F3-2 with  $\Gamma^0(A_0, A_1) = *uuv$  and  $\Gamma^0(A_2) = 3 * v$  as H<sup>2</sup> groups. •  ${}^{33}\Gamma_2(4u, 2v) = (r, r_2, m_3 - 1 = r_2^2 = m_3^2 = r^3 = (m_3r_2m_3r_2)^u = (m_3rm_3r^{-1})^v = (r_2r)^2, \ 2u \neq v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < 1, \ \frac{2}{3} + \frac{1}{v} < 1) = {}^{3}_{3}\Gamma_2(4u, 2v), \text{ with 2 extensions both for vertices } A_0, A_1 \text{ and } A_2$ (1) for  $A_0, A_1: (\bar{m}_0, \bar{m}_1 - 1 = \bar{m}_0^2 = \bar{m}_1^2 \stackrel{1}{=} \bar{m}_1r\bar{m}_0r^{-1} \stackrel{2}{=} \bar{m}_0r_2\bar{m}_1r_2 \stackrel{3}{=} (\bar{m}_0m_3)^2 \stackrel{4}{=} (\bar{m}_1m_3)^2)$ 

(1) for  $A_2$ :  $(\bar{m}_2 - 1 = \bar{m}_2^2 \stackrel{5}{=} \bar{m}_2 r \bar{m}_2 r^{-1} \stackrel{6}{=} (\bar{m}_2 m_3)^2)$ 

(2) for  $A_0$ ,  $A_1$ :  $(\bar{s} - 1 \stackrel{1}{=} (sr)^2 \stackrel{2}{=} (sr_2)^2 \stackrel{3}{=} \bar{s}m_3\bar{s}^{-1}m_3)$ (2) for  $A_2$ :  $(\bar{h}_2 - 1 = \bar{h}_2^2 \stackrel{5}{=} (\bar{h}_2r)^2 \stackrel{7}{=} \bar{h}_2m_3\bar{h}_2m_3)$ .

The pictures are given in Fig. 18: F3-3 with  $\Gamma^{0}(A_{0}, A_{1}) = 2 * uv$  and  $\Gamma^{0}(A_{2}) = 3 * v$  as  $\mathbb{H}^{2}$  groups. •  ${}^{33}\Gamma_{3}(4u, v) = (r, m_{2}, r_{3} - 1 = m_{2}^{2} = r_{3}^{2} = r^{3} = (m_{2}r_{3}m_{2}r_{3})^{u} = (r_{3}r)^{v} = m_{2}rm_{2}r^{-1}, 4u \neq v \in \mathbb{N}, \frac{1}{u} + \frac{2}{v} < 1, \frac{2}{3} + \frac{2}{v} < 1$ 1)  $= {}^{3}_{3}\Gamma_{3}(4u, v)$ , with 2 extensions both for vertices  $A_{0}, A_{1}$  and  $A_{2}$ (1) for  $A_{0}, A_{1}$ :  $(\bar{m}_{0}, \bar{m}_{1} - 1 = \bar{m}_{0}^{2} = \bar{m}_{1}^{2} = \bar{m}_{1}r\bar{m}_{0}r^{-1} \stackrel{2}{=} (\bar{m}_{0}m_{2})^{2} \stackrel{3}{=} \bar{m}_{0}r_{3}\bar{m}_{1}r_{3} \stackrel{4}{=} (\bar{m}_{1}m_{2})^{2})$ (1) for  $A_{2}$ :  $(\bar{m}_{2} - 1 = \bar{m}_{2}^{2} \stackrel{5}{=} \bar{m}_{2}r\bar{m}_{2}r^{-1} \stackrel{6}{=} \bar{m}_{2}r_{3}\bar{m}_{2}r_{3})$ (2) for  $A_{0}, A_{1}$ :  $(\bar{s} - 1 \stackrel{1}{=} (sr)^{2} \stackrel{2}{=} \bar{s}m_{2}\bar{s}^{-1}m_{2} \stackrel{3}{=} (\bar{s}r_{3})^{2})$ (2) for  $A_{2}$ :  $(\bar{h}_{2} - 1 = \bar{h}_{2}^{2} \stackrel{5}{=} (\bar{h}_{2}r)^{2} \stackrel{6}{=} (\bar{h}_{2}r_{3})^{2})$ .



Figure 8: F2-1 to  ${}^{2*2}\Gamma(u, 2v)$  and F2-2 to  ${}^{*22}\Gamma_2(2u, 2v)$ 



Figure 9: F2-3 to  ${}^{2\times}\Gamma_1(2u, 2v)$  and F2-4 to  ${}^{2\times}\Gamma_2(u, 2v)$ 



Figure 10: F2-5 to  $^{222}\Gamma_2(u, 2v)$  and F2-6<sub>1</sub> to  $^{22}\Gamma_7(2u, 4v)$ 



Figure 11: F2-6<sub>2</sub> to  ${}^{22}\Gamma_7(2u, 4v)$  simpler version



Figure 12: F2-7 to  ${}^{22}\Gamma_8(u, 4v)$ , simpler version is not pictured



Figure 13: F2-8<sub>1</sub> to  ${}^{22}\Gamma_9(2u, 2v)$  first version



Figure 14: F2-8<sub>2</sub> to  ${}^{22}\Gamma_9(2u, 2v)$  second simpler version



 $\Gamma_{58}(2u, 4v)$  fundamental 8 extensions

Figure 15: F2-8+ for  $\Gamma_{58}(2u, 4v)$  the 8 fundamental extensions (repetition)

 $\Gamma_{s8}(2u, 4v)$  fundamental 8 extensions cont.



(8) ext. by half turns at 2,3,5,6 to  $\tilde{s}_1, \tilde{s}_2$ 

Figure 16: F2-8+ for  $\Gamma_{58}(2u, 2v)$ , important topological phenomena



Figure 17: F3-1 to  $^{*33}\Gamma(2u, v)$  and F3-2 to  $^{33}\Gamma_1(2u, 2v)$ 



Figure 18: F3-3 to  ${}^{33}\Gamma_2(4u, 2v)$  and F3-4 to  ${}^{33}\Gamma_3(4u, v)$ 



Figure 19: F3-5 to  ${}^{33}\Gamma_4(2u, v)$ 

The pictures are given in Fig. 18: F3-4 with  $\Gamma^0(A_0, A_1) = v * u$  and  $\Gamma^0(A_2) = 23v$  as  $\mathbb{H}^2$  groups.

•  ${}^{33}\Gamma_4(2u,v) = (r,r_2,r_3-1=r_2^2=r_3^2=r^3=(r_2r_3)^u = (r_3r)^v = (rr_2)^2, \ 2u \neq v \in \mathbb{N}, \ \frac{2}{u} + \frac{2}{v} < 1, \ \frac{2}{3} + \frac{2}{v} < 1) = {}^{3}_{3}\Gamma_4(2u,v)$ , with 2 extensions both for vertices  $A_0, A_1$  and  $A_2$ (1) for  $A_0, A_1$ :  $(\bar{m}_0, \bar{m}_1 - 1 = \bar{m}_0^2 = \bar{m}_1^2 \stackrel{1}{=} \bar{m}_1 r \bar{m}_0 r^{-1} \stackrel{2}{=} \bar{m}_0 r_2 \bar{m}_1 r_2 \stackrel{3}{=} \bar{m}_0 r_3 \bar{m}_1 r_3)$ (1) for  $A_2$ :  $(\bar{m}_2 - 1 = \bar{m}_2^2 \stackrel{4}{=} \bar{m}_2 r \bar{m}_2 r^{-1} \stackrel{5}{=} \bar{m}_2 r_3 \bar{m}_2 r_3)$ (2) for  $A_0, A_1$ :  $(\bar{s} - 1 \stackrel{1}{=} (sr)^2 \stackrel{2}{=} (\bar{s}r_2)^2 \stackrel{3}{=} (\bar{s}r_3)^2)$ (2) for  $A_2$ :  $(\bar{h}_2 - 1 = \bar{h}_2^2 \stackrel{4}{=} (\bar{h}_2 r)^2 \stackrel{5}{=} (\bar{h}_2 r_3)^2)$ . The pictures are given in Fig. 19: F3-5 with  $\Gamma^0(A_0, A_1) = 2uv$  and  $\Gamma^0(A_2) = 23v$  as  $\mathbb{H}^2$  groups.

# 5. Family F4

The family F4 is characterized by  ${}^{*22}\Gamma_{1}(u, 2v, w) = (m_{0}, m_{1}, m_{2}, m_{3} - 1 = m_{0}^{2} = m_{1}^{2} = m_{2}^{2} = m_{3}^{2} = (m_{0}m_{1})^{w} =$   $= (m_{0}m_{2})^{2} = (m_{0}m_{3})^{2} = (m_{1}m_{2})^{2} = (m_{1}m_{3})^{v} = (m_{2}m_{3})^{u}; \ u \neq w, v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < \frac{1}{2}, \ \frac{1}{v} + \frac{1}{w} < \frac{1}{2}) =$   ${}^{mm_{4}^{2}}\Gamma_{1}(u, 2v, w), \text{ with trivial extensions both for vertices } A_{0} \text{ and } A_{2}$ (1) for  $A_{0}: (\bar{m}_{0} - 1 = \bar{m}_{0}^{2} \stackrel{1}{=} (\bar{m}_{0}m_{1})^{2} \stackrel{2}{=} (\bar{m}_{0}m_{2})^{2} \stackrel{3}{=} (\bar{m}_{0}m_{3})^{2}).$ (1) for  $A_{2}: (\bar{m}_{2} - 1 = \bar{m}_{2}^{2} \stackrel{4}{=} (\bar{m}_{2}m_{0})^{2} \stackrel{5}{=} (\bar{m}_{2}m_{1})^{2} \stackrel{6}{=} (\bar{m}_{2}m_{3})^{2}).$ The pictures are given in Fig. 20: F4-1 with  $\Gamma^{0}(A_{0}) = *2uv$  and  $\Gamma^{0}(A_{2}) = *2vw$  as  $\mathbb{H}^{2}$  groups. **Other non-fundamental cases from F4**  ${}^{*}\Gamma_{2}(u, 4v, w) = (r, m_{0}, m_{1} - 1 = m_{0}^{2} = m_{1}^{2} = r^{u} = (m_{1}rm_{1}r^{-1})^{v} = (m_{0}m_{1})^{w} = m_{0}rm_{0}r^{-1}, \ u \neq w, v \in \mathbb{N}, \ \frac{2}{u} + \frac{1}{v} < 1, \ \frac{2}{w} + \frac{1}{v} < 1) = \ {}^{m}_{2}\Gamma_{2}(u, 4v, w), \text{ with 2 extensions both for vertices } A_{0} \text{ and } A_{2}, A_{3}$ 

(1) for  $A_0$ :  $(\bar{m}_0 - 1 = \bar{m}_0^2 \stackrel{1}{=} (\bar{m}_0 m_1)^2 \stackrel{2}{=} \bar{m}_0 r \bar{m}_0 r^{-1})$ 

1436

(1) for  $A_2$ ,  $A_3$ :  $(\bar{m}_2, \bar{m}_3 - 1 = \bar{m}_2^2 = \bar{m}_3^2 \stackrel{3}{=} (\bar{m}_2 m_0)^2 \stackrel{4}{=} (\bar{m}_2 m_1)^2 \stackrel{5}{=} \bar{m}_3 r \bar{m}_2 r^{-1} \stackrel{6}{=} (\bar{m}_3 m_0)^2 \stackrel{7}{=} (\bar{m}_3 m_1)^2)$ (2) for  $A_0$ :  $(\bar{h}_0 - 1 = \bar{h}_0^2 \stackrel{1}{=} (m_1 \bar{h}_0)^2 \stackrel{2}{=} (\bar{h}_0 r)^2)$ (2) for  $A_2$ ,  $A_3$ :  $(\bar{s} - 1 \stackrel{3}{=} m_0 \bar{s} m_0 \bar{s}^{-1} \stackrel{4}{=} m_1 \bar{s} m_1 \bar{s}^{-1} \stackrel{5}{=} (r\bar{s})^2)$ . The pictures are given in Fig. 20: F4-2 with  $\Gamma^{0}(A_{0}) = u * v$  and  $\Gamma^{0}(A_{2}, A_{3}) = *vww$  as  $\mathbb{H}^{2}$  groups. •  ${}^{*}\Gamma_{3}(2u, 4v, 2w) = (r_{1}, m_{0}, m_{2}, m_{3} - 1 = r_{1}^{2} = m_{0}^{2} = m_{2}^{2} = m_{3}^{2} = (m_{0}m_{2})^{2} = (m_{0}m_{3})^{2} = (m_{2}m_{3})^{u} = (m_{2}r_{1}m_{3}r_{1})^{v} = (m_{2}r_{1}m_{3}r_{1})^{v} = (m_{2}r_{1}m_{3}r_{1})^{v} = (m_{2}r_{1}m_{3}r_{1})^{v}$  $(m_0 r_1 m_0 r_1)^w$ ,  $u \neq w, v \in \mathbb{N}$ ,  $\frac{1}{u} + \frac{1}{v} < 1$ ,  $\frac{1}{v} + \frac{1}{w} < 1$ ) =  $\frac{m}{2} \Gamma_3(2u, 4v, 2w)$ , with 2 extensions both for vertices  $A_0$ and  $A_2$ ,  $A_3$ (1) for  $A_0$ :  $(\bar{m}_0 - 1 = \bar{m}_0^2 \stackrel{1}{=} (\bar{m}_0 r_1)^2 \stackrel{2}{=} (\bar{m}_0 m_2)^2 \stackrel{3}{=} (\bar{m}_0 m_3)^2)$ (1) for  $A_2$ ,  $A_3$ :  $(\bar{m}_2, \bar{m}_3 - 1 = \bar{m}_2^2 = \bar{m}_3^2 \stackrel{4}{=} (\bar{m}_2 m_0)^2 \stackrel{5}{=} \bar{m}_2 r_1 \bar{m}_3 r_1 \stackrel{6}{=} (\bar{m}_2 m_3)^2 \stackrel{7}{=} (\bar{m}_3 m_2)^2 \stackrel{8}{=} (\bar{m}_3 m_0)^2$ ) (2) for  $A_0$ :  $(\bar{h}_0 - 1 = \bar{h}_0^2 \stackrel{1}{=} (r_1 \bar{h}_0)^2 \stackrel{2}{=} m_2 \bar{h}_0 m_3 \bar{h}_0)$ (2) for  $A_2$ ,  $A_3$ :  $(\bar{s} - 1 \stackrel{4}{=} m_0 \bar{s} m_0 \bar{s}^{-1} \stackrel{5}{=} (r_1 \bar{s})^2 \stackrel{6}{=} m_3 \bar{s} m_2 \bar{s}^{-1})$ . The pictures are given in Fig. 21: F4-3 with  $\Gamma^0(A_0) = 2 * uv$  and  $\Gamma^0(A_2, A_3) = *2v2w$  as  $\mathbb{H}^2$  groups. •  ${}^{*}\Gamma_{5}(u, 2v, 2w) = (m_{0}, r_{1}, r - 1 = m_{0}^{2} = r_{1}^{2} = r^{u} = (r_{1}r)^{v} = (m_{0}r_{1}m_{0}r_{1})^{w} = m_{0}rm_{0}r^{-1}, \ u \neq 2w, v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < 0$  $\frac{1}{2}, \frac{2}{v} + \frac{1}{w} < 1$  =  $\frac{m}{2}\Gamma_5(u, 2v, 2w)$ , with 2 extensions both for vertices  $A_0$  and  $A_2, A_3$ (1) for  $A_0$ :  $(\bar{m}_0 - 1 = \bar{m}_0^2 \stackrel{1}{=} (\bar{m}_0 r_1)^2 \stackrel{2}{=} \bar{m}_0 r \bar{m}_0 r^{-1})$ (1) for  $A_2, A_3$ :  $(\bar{m}_2, \bar{m}_3 - 1 = \bar{m}_2^2 = \bar{m}_3^2 \stackrel{3}{=} \bar{m}_2 r_1 \bar{m}_3 r_1 \stackrel{4}{=} \bar{m}_3 r \bar{m}_2 r^{-1} \stackrel{5}{=} (\bar{m}_2 m_0)^2 \stackrel{6}{=} (\bar{m}_3 m_0)^2$ ) (2) for  $A_0$ :  $(\bar{h}_0 - 1 = \bar{h}_0^2 \stackrel{1}{=} (\bar{h}_0 r_1)^2 \stackrel{2}{=} (\bar{h}_0 r)^2)$ (2) for  $A_2$ ,  $A_3$ :  $(\bar{s} - 1 \stackrel{3}{=} (r_1 \bar{s})^2 \stackrel{4}{=} (r \bar{s})^2 \stackrel{5}{=} m_0 \bar{s} m_0 \bar{s}^{-1})$ . The pictures are given in Fig. 21: F4-4 with  $\Gamma^0(A_0) = 2uv$  and  $\Gamma^0(A_2, A_3) = v * w$  as  $\mathbb{H}^2$  groups.

### 5.1. Truncated half-simplices from F4

•  ${}^{22}\Gamma_3(u, 4v, 2w) = (m_1, r_3, h-1 = m_1^2 = r_3^2 = h^2 = (r_3h)^u = (m_1r_3hm_1hr_3)^v = (m_1hm_1h)^w, \ u \neq 2w, v \in \mathbb{N}, \ \frac{2}{u} + \frac{1}{v} < 0$ 1,  $\frac{1}{n} + \frac{1}{m} < 1$  =  $\frac{2}{2} \Gamma_3(u, 4v, 2w)$ , with 2 extensions both for vertices  $A_0, A_1$  and  $A_2, A_3$ (1) for  $A_0, A_1$ :  $(\bar{m}_0, \bar{m}_1 - 1 = \bar{m}_0^2 = \bar{m}_1^2 \stackrel{1}{=} (m_1 \bar{m}_0)^2 \stackrel{2}{=} \bar{m}_0 h \bar{m}_1 h \stackrel{3}{=} \bar{m}_0 r_3 \bar{m}_1 r_3)$ (1) for  $A_2$ ,  $A_3$ :  $(\bar{m}_2, \bar{m}_3 - 1 = \bar{m}_2^2 = \bar{m}_3^2 \stackrel{4}{=} (\bar{m}_2 r_3)^2 \stackrel{5}{=} \bar{m}_2 h \bar{m}_3 h \stackrel{6}{=} (\bar{m}_2 m_1)^2 \stackrel{7}{=} (\bar{m}_3 m_1)^2)$ (2) for  $A_0$ ,  $A_1$  by half-turn  $\bar{h}_0$ :  $(\bar{h}_0 - 1 = \bar{h}_0^2 \stackrel{1}{=} (m_1 \bar{h}_0)^2 \stackrel{3}{=} (\bar{h}_0 h r_3)^2)$ (2) for  $A_2$ ,  $A_3$  by half-turn  $\bar{h}_2$ :  $(\bar{h}_2 - 1 \stackrel{4}{=} (\bar{h}_2 r_3)^2 \stackrel{6}{=} m_1 \bar{h}_2 h m_1 h \bar{h}_2)$ . The pictures are given in Fig. 22: F4-5 with  $\Gamma^0(A_0, A_1) = u * v$  and  $\Gamma^0(A_2, A_3) = 2 * vw$  as  $\mathbb{H}^2$  groups. •  ${}^{22}\Gamma_4(u, 2v, w) = (r_1, r_3, h - 1 = r_1^2 = r_3^2 = h^2 = (r_3h)^u = (r_1r_3h)^v = (r_1h)^w, \ u \neq w, v \in \mathbb{N}, \ \frac{1}{u} + \frac{1}{v} < \frac{1}{2}, \ \frac{1}{v} + \frac{1}{v} < \frac{1}{2}, \ \frac{1}{v} + \frac{1}{v} < \frac{1}{v} < \frac{1}{v} + \frac{1}{v} < \frac{1}{v} < \frac{1}{v} + \frac{1}{v} < \frac{1}{v} < \frac{1}{v} + \frac{1}{v} + \frac{1}{v} < \frac{1}{v} + \frac{1}{v} + \frac{1}{v} < \frac{1}{v} + \frac{1}{v} +$ 

 $\frac{1}{2}$ ) =  $\frac{2}{2} \Gamma_4(u, 2v, w)$ , with 2 extensions both for vertices  $A_0, A_1$  and  $A_2, A_3$ 

(1) for 
$$A_0, A_1$$
:  $(\bar{m}_0, \bar{m}_1 - 1 = \bar{m}_0^2 = \bar{m}_1^2 \stackrel{1}{=} (\bar{m}_0 r_1)^2 \stackrel{2}{=} \bar{m}_0 h \bar{m}_1 h \stackrel{3}{=} \bar{m}_0 r_3 \bar{m}_1 r_3)$ 

(1) for 
$$A_2$$
,  $A_3$ :  $(\bar{m}_2, \bar{m}_3 - 1 = \bar{m}_2^2 = \bar{m}_3^2 \stackrel{4}{=} (\bar{m}_2 r_3)^2 \stackrel{5}{=} \bar{m}_2 h \bar{m}_3 h \stackrel{6}{=} \bar{m}_2 r_1 \bar{m}_3 r_1$   
(2) for  $A_0$ ,  $A_1$  by half-turn  $\bar{h}_0$ :  $(\bar{h}_0 - 1 = \bar{h}_0^2 \stackrel{1}{=} (\bar{h}_0 r_1)^2 \stackrel{3}{=} (\bar{h}_0 h r_3)^2$ )

(2) for  $A_2$ ,  $A_3$  by half-turn  $\bar{h}_2$ :  $(\bar{h}_2 - 1 = \bar{h}_2^2 \stackrel{4}{=} (\bar{h}_2 r_3)^2 \stackrel{6}{=} (\bar{h}_2 h r_1)^2)$ .

The pictures are given in Fig. 23: F4-6 with  $\Gamma^0(A_0, A_1) = 2uv$  and  $\Gamma^0(A_2, A_3) = 2vw$  as  $\mathbb{H}^2$  groups.



Figure 20: F4-1 to  ${}^{*22}\Gamma_1(u, 2v, w)$  and F4-2 to  ${}^*\Gamma_2(u, 4v, w)$ 



 $\Gamma_{3}(2u, 4v, 2w) 4 (=2 \times 2)$  extensions

Figure 21: F4-3 to  ${}^*\Gamma_3(2u, 4v, 2w)$  and F4-4 to  ${}^*\Gamma_5(u, 2v, 2w)$ 



Figure 22: F4-5 to  ${}^{22}\Gamma_3(u, 4v, 2w)$ 



Figure 23: F4-6 to  ${}^{22}\Gamma_4(u, 4v, w)$ 

### 6. On packings and coverings with congruent hyperballs in Family 1-4

In this section we consider congruent hyperball packings and coverings for families F1-F2 and congruent hyperball packings for families F3-F4. The base planes of the hyperballs are always the truncation planes of the trunc-simplices.

- 1. First we consider Family 1, with a simple truncated orthoscheme (e.g.  $CHLA_1A_2A_3$  in Fig. 24.a and Fig. 1:F1-1) with Schläfli symbols  $\{\overline{u}, \overline{v}, \overline{w}\} = \{3, 3, u\} \ 6 < u \in \mathbb{N}$ ). This simple truncated orthoscheme can be derived also by truncation from orthoscheme  $A_0A_1A_2A_3 = b^0b^1b^2b^3$  with outer essential vertex  $A_0$ . The truncating plane  $a_0(a_0) = CLH$  is the polar plane of  $A_0$ , that is the (ultraparallel to  $b^0$ ) base plane of hyperball  $\mathcal{H}_0^h$  with height  $h_p$  for packing or  $h_c$  for covering.
- 2. For families F2-F4 (see Fig. 24.b and Fig. 8: F2-1, Fig. 17: F3-1, Fig. 20: F4-1, respectively) we consider a tiling  $\mathcal{T}(O(\overline{u}, \overline{v}, \overline{w}))$  with Schläfli symbol  $\{\overline{u}, \overline{v}, \overline{w}\}, (\frac{1}{\overline{u}} + \frac{1}{\overline{v}} < \frac{1}{2}, \frac{1}{\overline{v}} + \frac{1}{\overline{w}} < \frac{1}{2}, \overline{u}, \overline{v}, \overline{w} \in \mathbb{N})$  whose fundamental domain is doubly truncated orthoscheme  $O(\overline{u}, \overline{v}, \overline{w}) = CHLA_1A_2EJQ$  in Fig. 24.b.

Let a truncated orthoscheme  $O(\overline{u}, \overline{v}, \overline{w}) \subset \mathbf{H}^3$  be a tile from the above tiling. This can be derived also by truncation from  $A_0A_1A_2A_3 = b^0b^1b^2b^3$  with essential outer vertices  $A_0$  and  $A_3$ . The truncating planes  $a_0(a_0) = CLH$  and  $a_3(a_3) = JEQ$  are the polar planes of  $A_0$  and  $A_3$ , respectively, that *will be (ultraparallel to b<sup>0</sup> and b<sup>3</sup>) base planes of hyperballs*  $\mathcal{H}_i^h$  with height h (i = 0, 3). The distance between the two base planes  $d(a_0(a_0), a_3(a_3)) = d(H, J)$  (d is the hyperbolic distance function) will be double of the height of packing hyperball at most.



Figure 24: Simple and double truncated complete orthoschemes (without the absolute) in standard coordinate simplex  $b^0b^1b^2b^3 = A_0A_1A_2A_3$ ,  $(A_ib^j) = \delta_i^j$  with parameters  $\beta^{01} = \frac{\pi}{u}$ ,  $\beta^{12} = \frac{\pi}{u}$ ,  $\beta^{23} = \frac{\pi}{u}$ . Sketchy pictures also with half-turn axis  $F_{03}F_{12}$  in case  $\overline{u} = \overline{w}$ 

The Coxeter-Schläfli symbol of the complete orthoscheme tiling  $\mathcal{T}$  is generated by reflections in the planes  $b^i$  ( $i \in \{0, 1, 2, 3\}$ ) and  $a_j$  (j = 0, 3) of a complete orthoscheme O. To every scheme there is a corresponding symmetric  $4 \times 4$  matrix ( $b^{ij}$ ) where  $b^{ii} = 1$  and, for  $i \neq j \in \{0, 1, 2, 3\}$ ,  $b^{ij}$  equals to  $\cos(\pi - \beta^{ij}) = -\cos\beta^{ij}$  for all dihedral angles  $\beta^{ij}$  between the faces  $b^i, b^j$  of O.

For example, in formula (6.1) we see the Coxeter-Schläfli matrix with parameters  $(\overline{u}, \overline{v}, \overline{w})$ , i.e.  $\beta^{01} = \frac{\pi}{\overline{u}}$ ,  $\beta^{12} = \frac{\pi}{\overline{u}}$ ,  $\beta^{23} = \frac{\pi}{\overline{u}}$ .

$$(b^{ij}) = (\langle b^i, b^j \rangle) := \begin{pmatrix} 1 & -\cos\frac{\pi}{u} & 0 & 0 \\ -\cos\frac{\pi}{u} & 1 & -\cos\frac{\pi}{v} & 0 \\ 0 & -\cos\frac{\pi}{v} & 1 & -\cos\frac{\pi}{w} \\ 0 & 0 & -\cos\frac{\pi}{w} & 1 \end{pmatrix}.$$
(6.1)

This 3-dimensional complete (truncated or frustum) orthoscheme  $O = O(\overline{u}, \overline{v}, \overline{w})$  and its reflection group  $G_{\overline{u}\overline{v}\overline{w}}$  will be described in Fig. 24,1,8,17,20 and by the symmetric Coxeter-Schläfli matrix ( $b^{ij}$ ) in formula

(6.1), furthermore by its inverse matrix  $(a_{ij})$  in formula (6.2).

$$(a_{ij}) = (b^{ij})^{-1} = \langle \mathbf{a}_i, \mathbf{a}_j \rangle :=$$

$$= \frac{1}{B} \begin{pmatrix} \sin^2 \frac{\pi}{\overline{w}} - \cos^2 \frac{\pi}{\overline{v}} & \cos \frac{\pi}{\overline{u}} \sin^2 \frac{\pi}{\overline{w}} & \cos \frac{\pi}{\overline{u}} \cos \frac{\pi}{\overline{v}} & \cos \frac{\pi}{\overline{v}} \cos \frac{\pi}{\overline{v}} \cos \frac{\pi}{\overline{v}} \cos \frac{\pi}{\overline{v}} \\ \cos \frac{\pi}{\overline{u}} \sin^2 \frac{\pi}{\overline{w}} & \sin^2 \frac{\pi}{\overline{w}} & \cos \frac{\pi}{\overline{v}} \sin^2 \frac{\pi}{\overline{u}} \\ \cos \frac{\pi}{\overline{u}} \cos \frac{\pi}{\overline{v}} \cos \frac{\pi}{\overline{v}} \cos \frac{\pi}{\overline{v}} \cos \frac{\pi}{\overline{v}} \cos \frac{\pi}{\overline{v}} \sin^2 \frac{\pi}{\overline{u}} & \sin^2 \frac{\pi}{\overline{u}} - \cos^2 \frac{\pi}{\overline{v}} \end{pmatrix},$$

$$(6.2)$$

where

$$B = \det(b^{ij}) = \sin^2 \frac{\pi}{\overline{u}} \sin^2 \frac{\pi}{\overline{w}} - \cos^2 \frac{\pi}{\overline{v}} < 0, \text{ i.e. } \sin \frac{\pi}{\overline{u}} \sin \frac{\pi}{\overline{w}} - \cos \frac{\pi}{\overline{v}} < 0.$$

In the following the volume of  $O(\overline{u}, \overline{v}, \overline{w})$  is derived by the next Theorem of R. Kellerhals ([2], by the ideas of N. I. Lobachevsky):

**Theorem 3.** (R. Kellerhals) The volume of a three-dimensional hyperbolic complete orthoscheme  $O = O(\overline{u}, \overline{v}, \overline{w}) \subset \mathbb{H}^3$  can be expressed with the essential angles  $\beta^{01} = \frac{\pi}{\overline{u}}, \beta^{12} = \frac{\pi}{\overline{v}}, \beta^{23} = \frac{\pi}{\overline{w}}, (0 \leq \beta^{ij} \leq \frac{\pi}{2})$  in the following form:

$$\begin{aligned} \operatorname{Vol}(\mathcal{O}) &= \frac{1}{4} \{ \mathcal{L}(\beta^{01} + \theta) - \mathcal{L}(\beta^{01} - \theta) + \mathcal{L}(\frac{\pi}{2} + \beta^{12} - \theta) + \\ &+ \mathcal{L}(\frac{\pi}{2} - \beta^{12} - \theta) + \mathcal{L}(\beta^{23} + \theta) - \mathcal{L}(\beta^{23} - \theta) + 2\mathcal{L}(\frac{\pi}{2} - \theta) \}, \end{aligned}$$

where  $\theta \in (0, \frac{\pi}{2})$  is defined by:

$$\tan \theta = \frac{\sqrt{\cos^2 \beta^{12} - \sin^2 \beta^{01} \sin^2 \beta^{23}}}{\cos \beta^{01} \cos \beta^{23}}.$$

and where  $\mathcal{L}(x) := -\int_{0}^{x} \log |2 \sin t| dt$  denotes the Lobachevsky function (in J. Milnor's interpretation).

The hypersphere (or equidistant surface) is a quadratic surface at a constant distance from a plane (base plane) in both halfspaces. The infinite body of the hypersphere, containing the base plane, is called *hyperball*.

The *half hyperball* (i.e., the part of the hyperball lying on one side of its base plane) with distance *h* to a base plane  $\beta$  is denoted by  $\mathcal{H}^h_+$ . The volume of the intersection of  $\mathcal{H}^h_+(\mathcal{A})$  and the right prism with base a 2-polygon  $\mathcal{A} \subset \beta$ , can be determined by the classical formula of J. Bolyai.

$$\operatorname{Vol}(\mathcal{H}^{h}_{+}(\mathcal{A})) = \frac{1}{4}\operatorname{Area}(\mathcal{A})\left[k\sinh\frac{2h}{k} + 2h\right].$$
(6.3)

The constant  $k = \sqrt{\frac{-1}{K}}$  is the natural length unit in  $\mathbb{H}^3$ , where *K* denotes the constant negative sectional curvature. In the following we may assume that k = 1. The distance *d* of two proper points *X*(**x**) and *Y*(**y**) is calculated by the formula

$$\cosh d = \frac{-\langle \mathbf{x}, \mathbf{y} \rangle}{\sqrt{\langle \mathbf{x}, \mathbf{x} \rangle \langle \mathbf{y}, \mathbf{y} \rangle}}.$$
(6.4)

# 6.1. Congruent hyperballs to a simply truncated orthoscheme, Family 1

The maximal series of family F1 is characterized by the simplex tiling whose group series is  ${}^{\bar{4}3m}_{24}\Gamma(u)$  with previous crystallographic notation from [10], and now with the new orbifold notation  ${}^{*233}\Gamma(u)$ . It has the trivial extension by the reflection  $a_0 = \bar{m}_0$ , as in Sect. 2.

The corresponding unique fundamental simplex with parameters  $(\overline{u}, \overline{v}, \overline{w}) = (3, 3, u)$  is determined by its Coxeter-Schläfli matrix which is derived by (6.1). Its determinant is  $B = \det(b^{ij}) = \frac{3}{4} \sin^2 \frac{\pi}{u} - \frac{1}{4} < 0$  now  $6 < u \in \mathbb{N}$ .

1443

So, the signature is (+ + + -) and the simplex is realizable in hyperbolic space  $\mathbb{H}^3$ . The inverse  $(a_{ij})$ of the Coxeter-Schläfli matrix  $(b^{ij})$  follows by (6.2) providing the distance metric (by the general theory of projective metrics [13]).

The optimal hyperball height is  $h_p = d(a_0, m_0) = d(C, A_1)$  for packings and  $h_c = d(H, A_3)$  is the optimal covering distance where

$$h_{p} = \cosh d(C, A_{1}) = \frac{-a_{00}a_{11} + a_{01}^{2}}{\sqrt{a_{00}a_{11}\left(-a_{01}^{2} + a_{00}a_{11}\right)}} = \sqrt{\frac{\frac{1}{4} - \frac{3}{4}\sin^{2}\frac{\pi}{u}}{\frac{1}{4} - \sin^{2}\frac{\pi}{u}}},$$

$$h_{c} = \cosh d(H, A_{3}) = \frac{-a_{00}a_{33} + a_{03}^{2}}{\sqrt{a_{00}a_{33}\left(-a_{03}^{2} + a_{00}a_{33}\right)}} = \sqrt{\frac{\frac{1}{2}\left(\frac{1}{4} - \sin^{2}\frac{\pi}{u}\right) + \frac{1}{16}\cos^{2}\frac{\pi}{u}}{\frac{1}{2}\left(\frac{1}{4} - \sin^{2}\frac{\pi}{u}\right)}}$$
equivalent to  $\sinh d(C, A_{1}) = \frac{\frac{1}{2}\sin\frac{\pi}{u}}{\sqrt{\frac{1}{4} - \sin^{2}\frac{\pi}{u}}}$  and  $\sinh d(H, A_{3}) = \frac{\frac{1}{4}\cos\frac{\pi}{u}}{\sqrt{\frac{1}{2}\left(\frac{1}{4} - \sin^{2}\frac{\pi}{u}\right)}}.$ 

Then we get the densities  $\delta$  (packing),  $\Delta$  (covering), respectively. E.g.

$$\delta(O(3,3,u)) = \frac{\operatorname{Area}(\mathcal{A})\frac{1}{4}\left[\sinh 2h_p + 2h_p\right]}{\operatorname{Vol}(O)},$$

where Area( $\mathcal{A}$ ) =  $\left(\frac{\pi}{2} - \frac{\pi}{3} - \frac{\pi}{u}\right)$ . The Vol(O) can be calculated by Theorem 3.

In this case the maximal packing density is  $\approx 0.82251$  with u = 7 (see in Table 1p). The optimal covering density of Family 1 is  $\approx$  1.33093 for u = 7 in Table 1c.

<b>Table 1p</b> , $(3, 3, u)$ , $6 < u \in \mathbb{N}$					
и	h	Vol(O)	$\operatorname{Vol}(\mathcal{H}^h_+(\mathcal{A}))$	δ	
7	0.78871	0.08856	0.07284	0.82251	
8	0.56419	0.10721	0.08220	0.76673	
9	0.45320	0.11825	0.08474	0.71663	
÷	:	:	:	÷	
20	0.16397	0.14636	0.06064	0.41431	
÷	:	:	÷	:	
50	0.06325	0.15167	0.02918	0.19240	
÷	:	÷		:	
100	0.03147	0.15241	0.01549	0.10165	
$u \to \infty$	0	0.15266	0	0	
	Ta	<b>ble 1c,</b> (3, 3, 1	$u$ ) 6 < $u \in \mathbb{N}$		
и	h	Vol(O)	$\operatorname{Vol}(\mathcal{H}^h_+(\mathcal{A}))$	Δ	
7	1.06739	0.08856	0.11787	1.33093	
8	0.89198	0.10721	0.15304	1.42747	
9	0.81696	0.11825	0.17882	1.51225	
÷	:	:	:	÷	
20	0.68136	0.14636	0.29213	1.99596	
:	:	:	:	:	
50	0.66193	0.15167	0.35361	2.33146	
:	:	:	:	:	

0.65934 0.15241

0.15266

100  $u \rightarrow \infty$  0.65848 0.37580

0.39911

2.46566

2.61438

By completing Tables 1p and 1c we can read the optimal hyperball packings and coverings with their densities for other extensions in F1 as well. E.g. to  ${}^{233}\Gamma(2u)$  we get these for 2u = 8 with  $\delta = \approx 0.76673$ ,  $\Delta \approx 1.42747$ . To  ${}^{2\times}\Gamma_3(3u)$  we get these for 3u = 9 with  $\delta \approx 0.71663$ ,  $\Delta \approx 1.51225$ . To  ${}^{33}\Gamma_5(4u)$  we get these for 4u = 8 again with  $\delta \approx 0.76673$ ,  $\Delta \approx 1.42747$ .

The argument is by our convention that the parameters of groups above in parentheses are just the  $m^{23}$ entries of the *D*-matrix function,  $2m^{23}$  is the number of subsimplices around simplex edges. Mostow's rigidity principle for co-compact hyperbolic space groups are also used (see e.g. [27]).

### 6.2. On hyperball packings and coverings in a doubly truncated orthoscheme, Family 2

In this case we consider a doubly truncated orthoscheme with Schläfli symbol  $\{\overline{u}, \overline{v}, \overline{w} = \overline{u}\}, (\frac{1}{\overline{u}} + \frac{1}{\overline{v}} < \frac{1}{2})$  $3 \leq \overline{u}, \overline{v} \in \mathbb{N}$ ) whose fundamental domain is *CHLA*<sub>1</sub>*A*<sub>2</sub>*EJQ* in Fig. 24.b. Its Coxeter Schläfli matrix  $(b^{ij})$  and its inverse  $(a_{ij})$  can be derived by formulas (6.1) and (6.2).

Volume formula of half hyperball  $H^h_+(\mathcal{R}_i)$  (i = 0, 3) of height h can be calculated by (6.3) where Area( $\mathcal{R}_0$ ) = Area $(\mathcal{A}_3) = \pi \left(\frac{1}{2} - \frac{1}{\bar{u}} - \frac{1}{\bar{v}}\right)$  according to the above orthoscheme.

Both polar planes assigne hyperspheres that are congruent with each other therefore the heights  $h^0 = h^3$ of optimal hyperballs  $\mathcal{H}_{i}^{h^{t}}$  (*i* = 0, 3) can be computed by Fig. 24.b.

The maximal height for optimal congruent hyperball packing is  $h_v = \min \{d(A_2, a_3), d(F_{03}, a_3)\}$ .

In the above doubly truncated orthoscheme  $d(A_2, a_3) = d(A_2, Q \sim a_{23}\mathbf{a}_3 - a_{33}\mathbf{a}_2), d(F_{03}, a_3) = d(\mathbf{a}_0 + \mathbf{a}_3)$  $\mathbf{a}_{3}, a_{33}\mathbf{a}_{0} - a_{03}\mathbf{a}_{3} \sim I$ , (as in [13]). Thus,

$$\cosh\left(d(A_2, a_3)\right) = \frac{-a_{33}a_{22} + a_{23}^2}{\sqrt{a_{22}\left(a_{33}^2 a_{23} - a_{23}^2 a_{33}\right)}} = \sqrt{1 - \frac{a_{23}^2}{a_{22}a_{33}}},$$

and

$$\cosh\left(d(F_{03}, a_3)\right) = \sqrt{\frac{a_{33} - a_{03}}{2a_{33}}} = \sqrt{\frac{1}{2} \left(1 + \frac{\cos^2 \frac{\pi}{\bar{u}} \cos \frac{\pi}{\bar{v}}}{\cos^2 \frac{\pi}{\bar{v}} - \sin^2 \frac{\pi}{\bar{u}}}\right)}$$

The density of optimal hyperball packing will be determined, by the following formula:

,

$$\delta(O(\overline{u},\overline{v},\overline{u})) = \frac{\operatorname{Area}(\mathcal{A})(\sinh 2h_p + 2h_p)}{2\operatorname{Vol}(O)} = \frac{\pi\left(\frac{1}{2} - \frac{1}{\overline{u}} - \frac{1}{\overline{v}}\right)\left(\sinh 2h_p + 2h_p\right)}{2\operatorname{Vol}(O)}.$$
(6.5)

### For the minimal covering hyperball height:

 $h_c = d(F_{12}, a_0) = d(F_{12}, a_3) = d(\mathbf{a}_1 + \mathbf{a}_2, \mathbf{a}_1 + \mathbf{a}_2 + c\mathbf{a}_3),$ 

where  $a_{13} + a_{23} + ca_{33} = 0$ , i.e.  $c = \frac{-(a_{13} + a_{23})}{a_{33}}$ . Thus, by  $a_{11} = a_{22}$ , we get

$$\cosh\left(d(F_{12}, a_3) = \frac{-\left(2\left(a_{22} + a_{12}\right) - \frac{1}{a_{33}}\left(a_{13} + a_{23}\right)^2\right)}{\sqrt{2\left(a_{22} + a_{12}\right)\left(2a_{22} + 2a_{12} - \frac{1}{a_{33}}\left(a_{13} + a_{23}\right)^2\right)}}$$

The covering density  $\Delta(O(\overline{u}, \overline{v}, \overline{u}))$  can be defined similarly to the packing density (see (6.5)).

 $Vol(O(\bar{u}, \bar{v}, \bar{u}))$  can be calculated by Theorem 3. The maximal volume sum of the hyperball pieces lying in  $O(\overline{u}, \overline{v}, \overline{u})$  can be computed by the formulas (6.1-4) and by the above described computation method for each given possible parameters  $\overline{u}, \overline{v}, \overline{u}$ . Therefore, the maximal packing density and the thinnest covering density of the congruent hyperball packings and coverings can be computed for each possible parameters.

We see the optimal packing density for F2 is  $\approx 0,81335$  with  $\overline{u} = \overline{w} = 7, \overline{v} = 3$  in Table 2p and the top minimal covering density is  $\approx 1,26869$  with  $\overline{u} = \overline{w} = 7$ , (as before)  $\overline{v} = 3$  in Table 2c. See Fig. 8: F2–1 and Fig. 25.

<b>Table 2p,</b> $(\overline{u}, \overline{v}, \overline{u})$						
$\{\overline{u}, \overline{v}, \overline{u}\}$	$h_p$	$Vol(O(\overline{u}, \overline{v}, \overline{u}))$	$\sum_{i=0,3} Vol(\mathcal{H}^h(\mathcal{A}_i))$	$\delta(O(\overline{u}, \overline{v}, \overline{u}))$		
{7,3,7}	1.23469	0.38325	0.31172	0.81335		
{6,4,6}	0.69217;	0.55557	0.42610	0.76696		
{8,3,8}	0.94946	0.44383	0.33794	0.76143		
{8,4,8}	0.56419	0.64328	0.49322	0.76673		
{5,4,5}	0.88055	0.46190	0.36007	0.77955		
{4,5,4}	0.80846	0.43062	0.31702	0.73620		
{4,6,4}	0.57311	0.50192	0.33516	0.66775		
{3,7,3}	0.98399	0.27899	0.20481	0.73411		
Table 2c, $(\overline{u}, \overline{v}, \overline{u})$						
		Table 2	$(\overline{u},\overline{v},\overline{u})$			
$\overline{\{\overline{u},\overline{v},\overline{u}\}}$	h <sub>c</sub>	Table 2 $Vol(O(\overline{u}, \overline{v}, \overline{u}))$	$\frac{\sum_{i=0,3} Vol(\mathcal{H}^h(\mathcal{A}_i))}{\sum_{i=0,3} Vol(\mathcal{H}^h(\mathcal{A}_i))}$	$\delta(O(\overline{u},\overline{v},\overline{u}))$		
$\overline{\{\overline{u},\overline{v},\overline{u}\}}$ $\overline{\{7,3,7\}}$	<i>h</i> <sub>c</sub> 1.49903	Table 20 $Vol(O(\overline{u}, \overline{v}, \overline{u}))$ 0.38325	$\frac{\sum_{i=0,3} Vol(\mathcal{H}^h(\mathcal{A}_i))}{0.48607}$	$\frac{\delta(O(\overline{u},\overline{v},\overline{u}))}{1.26829}$		
$     {\overline{u}, \overline{v}, \overline{u}}     {\overline{7, 3, 7}}     {\overline{6, 4, 6}} $	$h_c$ 1.49903 1.01481	Table 20 $Vol(O(\overline{u}, \overline{v}, \overline{u}))$ 0.38325           0.55557         0.55557	$\frac{c, (\overline{u}, \overline{v}, \overline{u})}{\sum_{i=0,3} Vol(\mathcal{H}^h(\mathcal{A}_i))}$ 0.48607 0.75523	$\frac{\delta(O(\overline{u},\overline{v},\overline{u}))}{1.26829}$ 1.35938		
	$\begin{array}{c} h_c \\ 1.49903 \\ 1.01481 \\ 1.26595 \end{array}$	Table 2 $Vol(O(\overline{u}, \overline{v}, \overline{u}))$ 0.38325           0.55557           0.44383	$\frac{\sum_{i=0,3} Vol(\mathcal{H}^{li}(\mathcal{A}_i))}{0.48607} \\ 0.75523 \\ 0.57470$	$\frac{\delta(O(\bar{u}, \bar{v}, \bar{u}))}{1.26829}$ 1.35938 1.29487		
$     \begin{array}{c} \overline{\{\overline{u},\overline{v},\overline{u}\}} \\ \overline{\{7,3,7\}} \\ \overline{\{6,4,6\}} \\ \overline{\{8,3,8\}} \\ \overline{\{5,4,5\}} \end{array} $	$\begin{array}{c} h_c \\ 1.49903 \\ 1.01481 \\ 1.26595 \\ 1.19095 \end{array}$	Table 2. $Vol(O(\overline{u}, \overline{v}, \overline{u}))$ 0.38325         0.55557         0.44383         0.46190	$ \begin{array}{c} \sum_{i=0,3} Vol(\mathcal{H}^{i}(\mathcal{A}_{i})) \\ \hline 0.48607 \\ 0.75523 \\ 0.57470 \\ \hline 0.60856 \end{array} $	$\frac{\delta(O(\bar{u}, \bar{v}, \bar{u}))}{1.26829}$ 1.35938 1.29487 1.31751		
	$\begin{array}{c} h_c \\ 1.49903 \\ 1.01481 \\ 1.26595 \\ 1.19095 \\ 0.89198 \end{array}$	Table 2.           Vol(O(ū, v, v))           0.38325           0.55557           0.44383           0.46190           0.64328	$ \begin{array}{c} \sum_{i=0,3} Vol(\mathcal{H}^{i}(\mathcal{A}_{i})) \\ \hline 0.48607 \\ 0.75523 \\ 0.57470 \\ \hline 0.60856 \\ 0.91826 \end{array} $	$\frac{\delta(O(\bar{u}, \bar{v}, \bar{u}))}{1.26829}$ 1.35938 1.29487 1.31751 1.42747		
	$\begin{array}{c} h_c \\ 1.49903 \\ 1.01481 \\ 1.26595 \\ 1.19095 \\ 0.89198 \\ 1.16974 \end{array}$	Table 20           Vol(O(\overline{n}, \overline{n}, \overline{n}))           0.38325           0.55557           0.44383           0.46190           0.64328           0.43062	$ \begin{array}{c} \overline{z}, (\overline{u}, \overline{v}, \overline{u}) \\ \hline \Sigma_{i=0,3} \ Vol(\mathcal{H}^{h}(\mathcal{A}_{i})) \\ 0.48607 \\ 0.75523 \\ 0.57470 \\ 0.60856 \\ 0.91826 \\ 0.58741 \end{array} $	$\frac{\delta(O(\bar{u},\bar{v},\bar{u}))}{1.26829}$ 1.35938 1.29487 1.31751 1.42747 1.36411		
	<i>h</i> <sub>c</sub> 1.49903 1.01481 1.26595 1.19095 0.89198 1.16974 0.99583	Table 20           Vol(O(i, v, v, i))           0.38325           0.55557           0.44383           0.46190           0.64328           0.43062           0.50192	$ \begin{array}{c} \overline{z}, (\overline{u}, \overline{v}, \overline{u}) \\ \hline \Sigma_{i=0,3} Vol(\mathcal{H}^{h}(\mathcal{A}_{i})) \\ 0.48607 \\ 0.75523 \\ 0.57470 \\ 0.60856 \\ 0.91826 \\ 0.58741 \\ 0.73137 \end{array} $	$\begin{array}{c} \hline \delta(O(\overline{u},\overline{v},\overline{u})) \\ \hline 1.26829 \\ \hline 1.35938 \\ \hline 1.29487 \\ \hline 1.31751 \\ \hline 1.42747 \\ \hline 1.36411 \\ \hline 1.45714 \end{array}$		

By completing Tables 2p and 2c we can determine the other densities in Family 2. E.g. to  ${}^{*22}\Gamma_2(2u, 2v)$ (i.e. u = 4, v = 3) we get  $\delta \approx 0.76143$ .  $\Delta \approx 1.29487$ . Or to  $\Gamma_{58}(2u, 4v)$  we get  $\{\overline{u}, \overline{v}, \overline{w}\} = \{6, 4, 6\}$  for optimal arrangement,  $\delta \approx 0.76696$ .  $\Delta \approx 1.35938$ .

### 6.3. Packings with congruent hyperballs in doubly truncated orthoscheme, Family 3-4

In this subsection we consider congruent hyperball packings to truncated orthoscheme tilings. Both polar planes assign hyperballs with equal height (see Fig. 24.b). It is clear, that the height of optimal hyperballs  $\mathcal{H}_i^{h^i}$  (*i* = 0, 3) is

$$h_p(\overline{u}, \overline{v}, \overline{w}) = \min\{d(H, J)/2, d(Q, A_2), d(C, A_1)\},\$$

where  $\overline{u}, \overline{v}, \overline{w}$  are integer parameters. In this case the volume sum of the hyperball pieces lying in the orthoscheme has to divide with the volume of trunc-orthoscheme  $O(\overline{u}, \overline{v}, \overline{w})$  as usual. Segments  $A_1C$ ,  $A_2Q$  and *JH* can be determined by the machinery of the projective metrics (see subsection 6.1 or [13]).

The volume of the orthoscheme  $O(\overline{u}, \overline{v}, \overline{w})$  can be determined by Theorem 3. We note here, that the role of the parameters u and w is symmetrical, therefore we can assume that  $\overline{u} > \overline{w}$  in F3 or  $\overline{u} < \overline{w}$  in F4.

<b>Table 3p</b> , $(\overline{u}, \overline{v}, \overline{w} = 3)$ , congruent hyperballs				
$\{\overline{u}, \overline{v}, \overline{w}\}$	h	$Vol(O(\overline{u}, \overline{v}, \overline{w}))$	$\sum_{i=0,3} Vol(\mathcal{H}^h(\mathcal{A}_i))$	$\delta^1(O(\overline{u},\overline{v},\overline{w}))$
{4,7,3}	0.59710	0.39274	0.27700	0.70529
{5,7,3}	0.41812;	0.43216	0.25203	0.58320
:		-	-	:
{50,7,3}	0.03492	0.49140	0.03962	0.08062
•	:	- - -	-	
{4,8,3}	0.56419	0.42885	0.32881	0.76673
{5,8,3}	0.67409	0.47536	0.26747	0.56266
•			:	
{50, 8, 3}	0.03405	0.52378	0.04245	0.08105
•		•	-	-
{4,9,3}	0.46841	0.45130	0.30800	0.68247
{5,9,3}	0.39083	0.48771	0.31589	0.64771
:	:	-	-	-
{50,9,3}	0.03348	0.54384	0.04466	0.08212

E. Molr	ıár et al./	Filomat	37:5 (	(2023),	1409–1448
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Table 4p, $(\overline{u}, \overline{v}, \overline{w})$ , congruent hyperballs				
$\{\overline{u},\overline{v},\overline{w}\}$	h	$Vol(O(\overline{u}, \overline{v}, \overline{w}))$	$\sum_{i=0,3} Vol(\mathcal{H}^h(\mathcal{A}_i))$	$\delta^1(O(\overline{u}, \overline{v}, \overline{w}))$
{7,3,8}	0.93100	0.41326	0.25726	0.62251
{7,3,9}	0.76734	0.43171	0.23355	0.54099
:	:	÷	:	:
{7,3,50}	0.11380	0.49016	0.06121	0.12488
•	-	:	-	
{8,3,9}	0.78366	0.46266	0.29474	0.63704
{8,3,10}	0.67409	0.47536	0.26747	0.56266
:	:	-	:	
{8,3,50}	0.11668	0.52248	0.06935	0.13274
:	:	:	:	
{5,4,6}	0.73969	0.50747	0.37287	0.73476
{5,4,7}	0.59326	0.53230	0.32974	0.61947
:	:		:	
{5,4,50}	0.07206	0.59291	0.06350	0.10710
· · · · ·			-	
{4,5,5}	0.69129	0.49789	0.38284	0.76893
{4,5,6}	0.53064	0.52971	0.33597	0.63426
:	:	:	:	
{4,5,50}	0.05502	0.59318	0.05710	0.096256
:	:		:	
{4,6,5}	0.50625	0.55992	0.37558	0.67078
{4,6,6}	0.48121	0.58850	0.40850	0.69414
	:	:		
$\{4, 6, 50\}$	0.05138	0.64697	0.06409	0.09906



Figure 25: The least dense congruent hyperball covering arrangement to parameters  $\{7, 3, 7\}$  with density  $\approx 1.26829$  and the corresponding Coxeter orthoscheme with half-turn axis in Euclidean B-C-K model. Compare also with Fig. 8: F2–1.

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